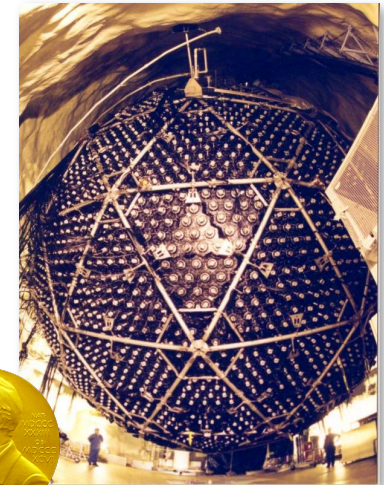
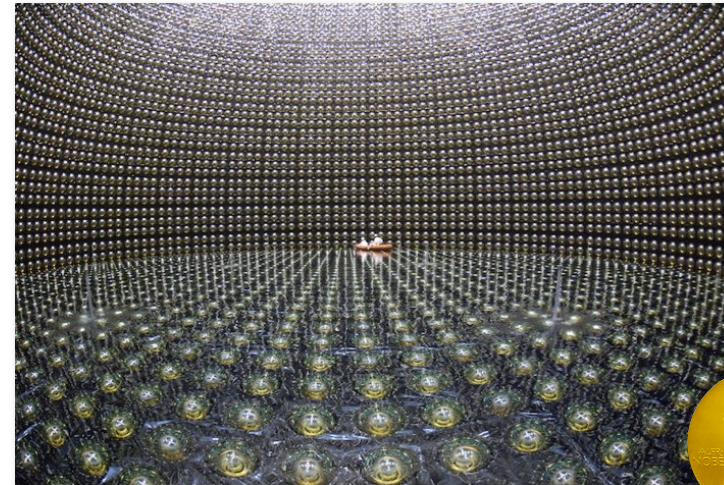


## III. Neutrino flavour oscillations



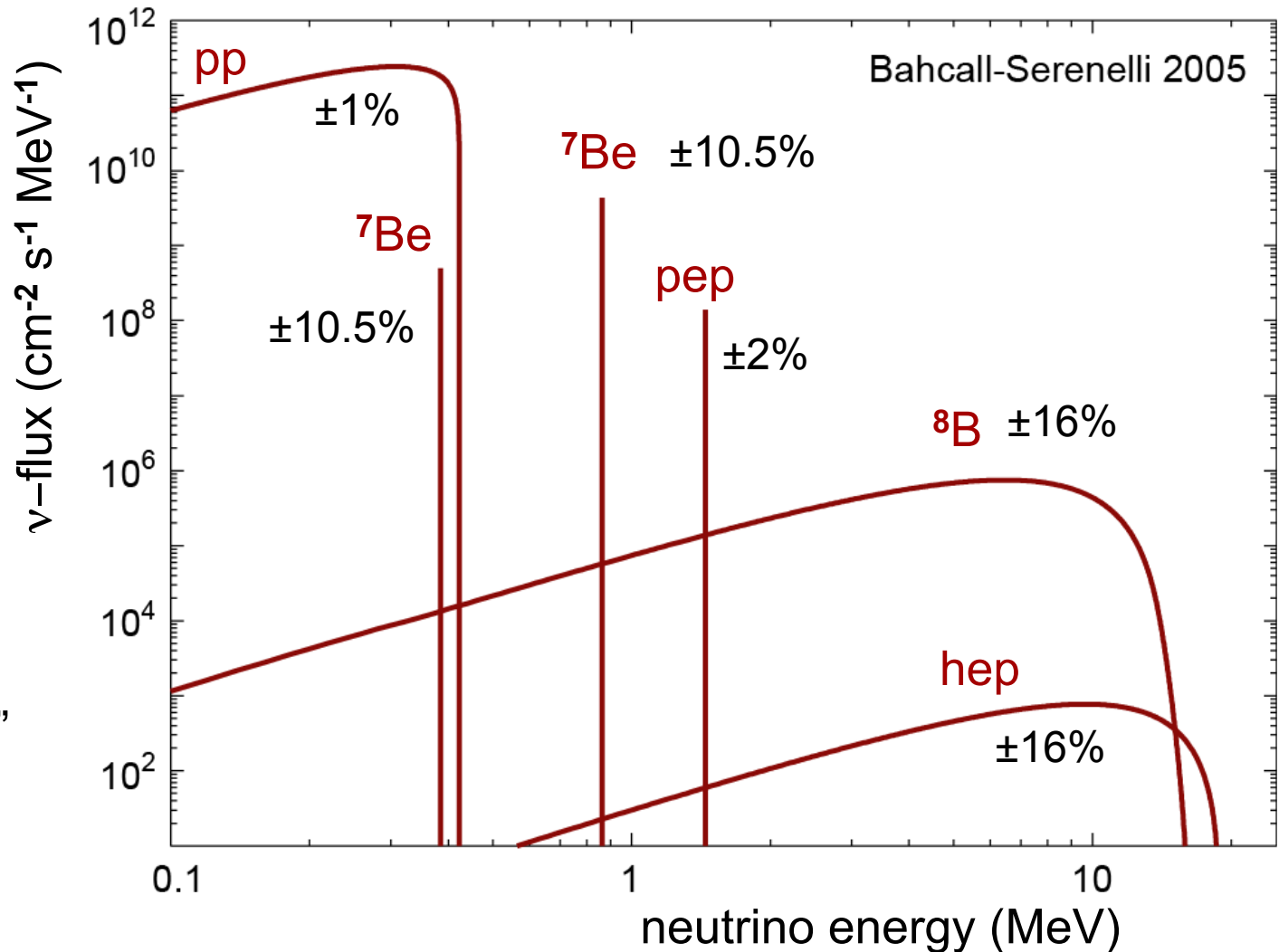
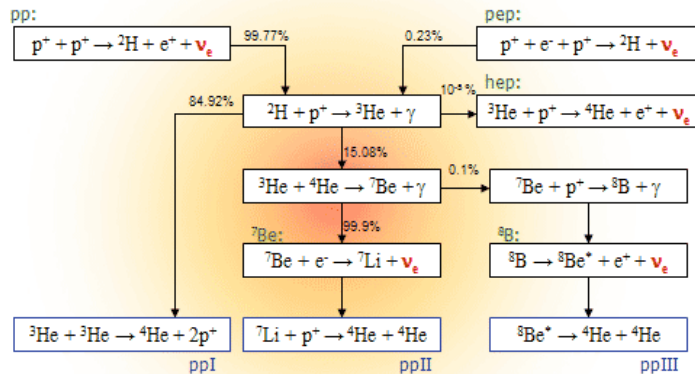
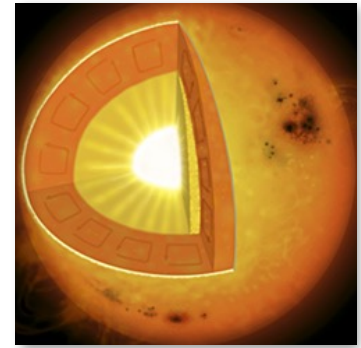
**Physics Nobel Prize 2002**  
to R. Davis Jr., M. Koshiba  
*“detection of cosmic [solar] neutrinos”*



**Physics Nobel Prize 2015**  
to A. B. McDonald and T. Kajita  
*“discovery of neutrino oscillations  
which show that neutrinos have a mass”*

# Neutrinos from the sun

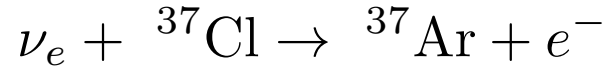
- Nuclear fusion in the solar core ( $T \approx 14.5 \times 10^6$  K)
- Only electron neutrinos are created
- Integral flux  $\approx 66$  billion  $\nu$  /cm<sup>2</sup> / s on Earth



Since 1960s:  
“Standard Solar Model”  
with detailed  $\nu$  flux  
calculations

# Experimental test of the solar model

low-threshold reaction



“radiochemical method”: extract noble gas argon and detect its decay ( $T_{1/2} = 35$  d)

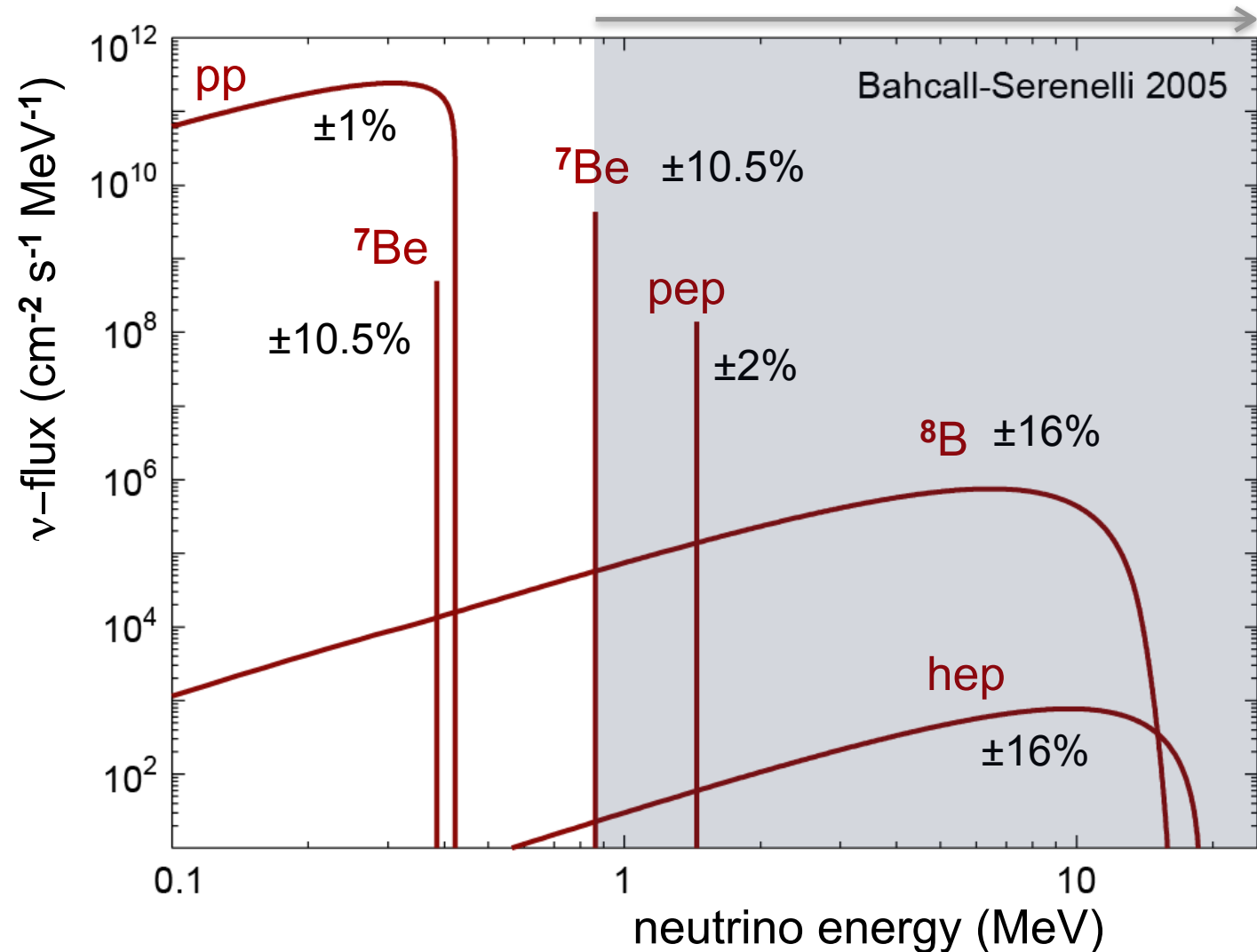
Ray Davis and John Bahcall  
at the Homestake mine, ca. 1964



615 tons of perchloroethylene  
buried in a gold mine



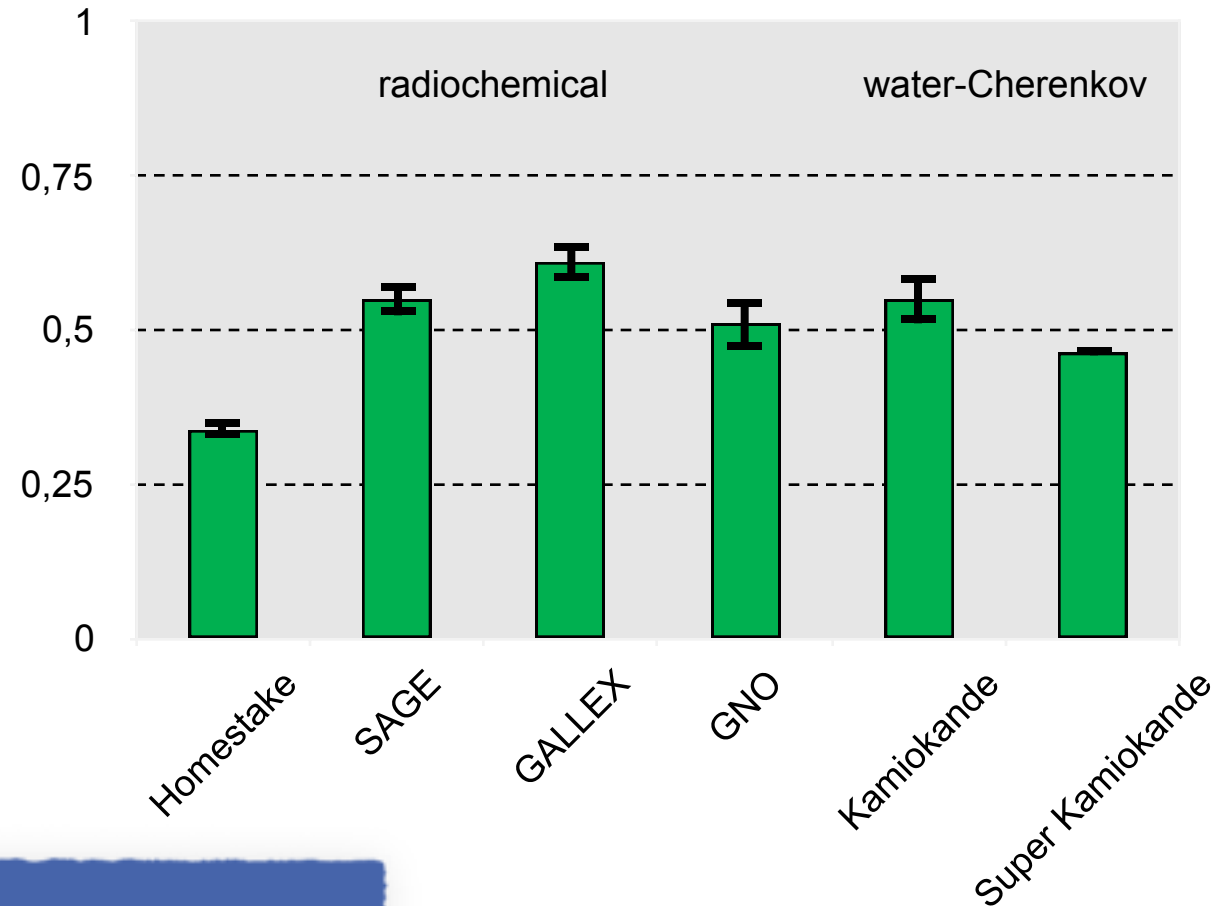
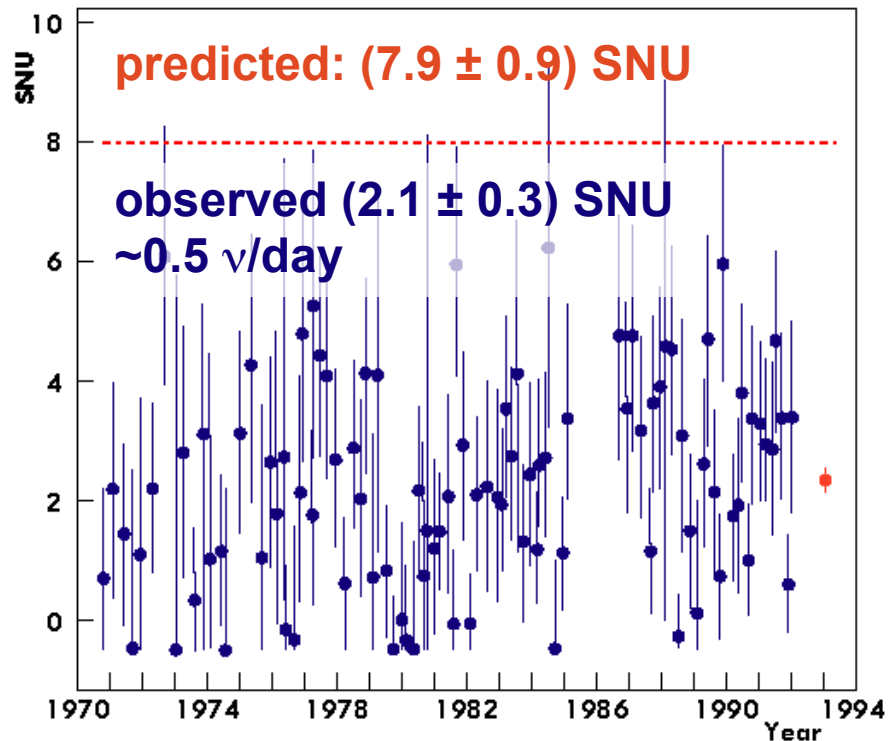
1 neutrino every 2 days





# The solar neutrino puzzle

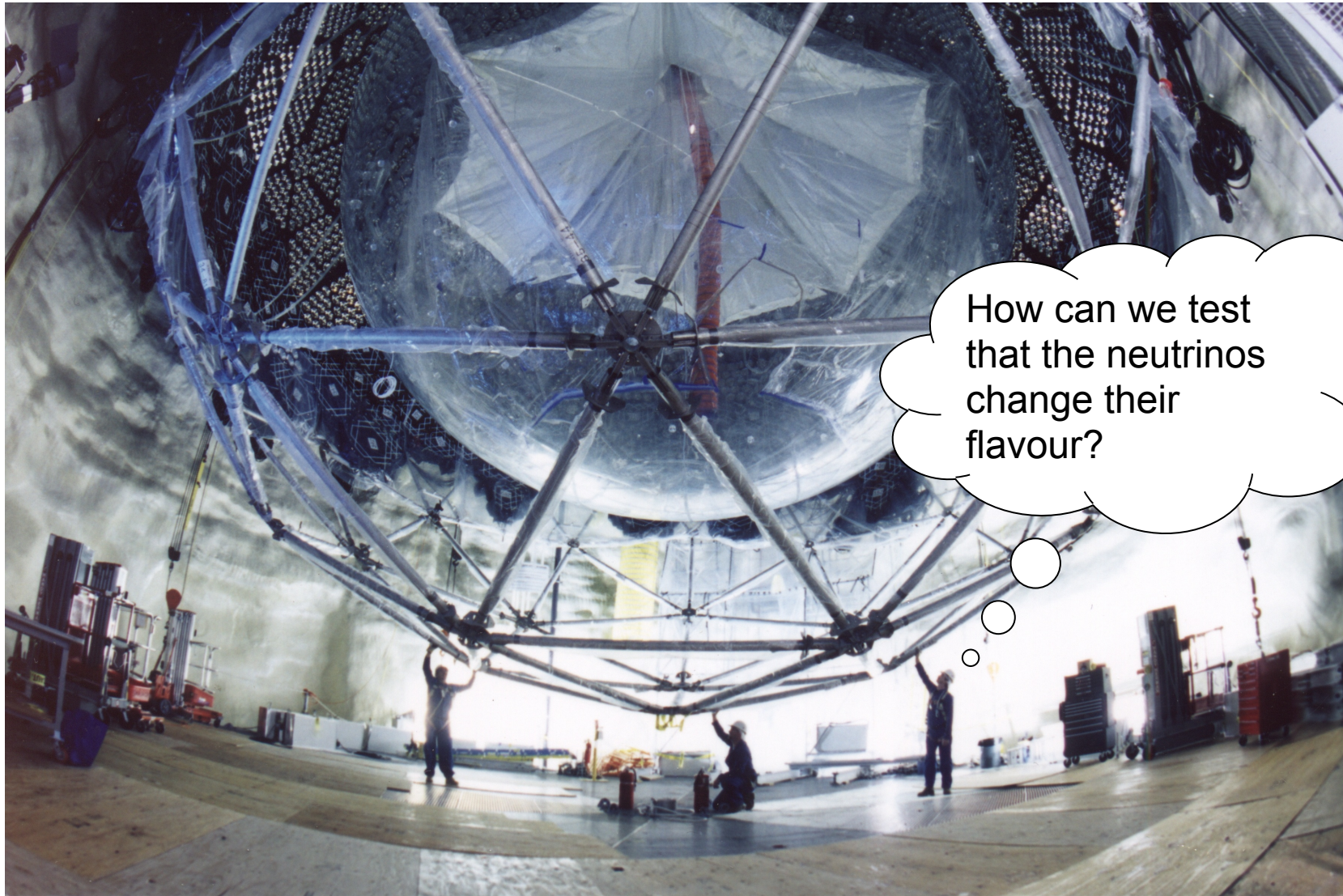
Too few neutrinos detected, consistently!



- Something wrong with all the experiments ?
- Something wrong with the solar model ?
- Something going on with the neutrinos ??



# SNO provides the answer to the problem

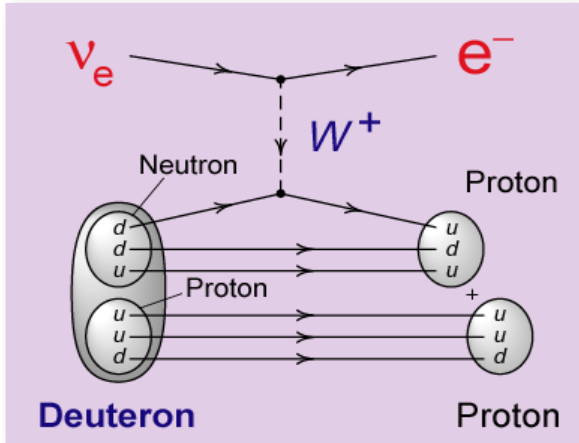


How can we test that the neutrinos change their flavour?

The **Sudbury Neutrino Observatory**, Creighton mine, Ontario/Canada (2100 m deep)  
1000 tons of heavy water ( $D_2O$ ) viewed by 9600 PMTs

# The SNO idea

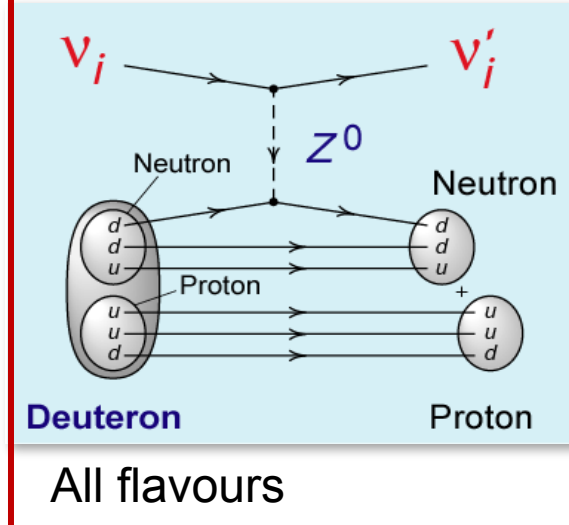
## Charged current (CC)



Only  $\nu_e$

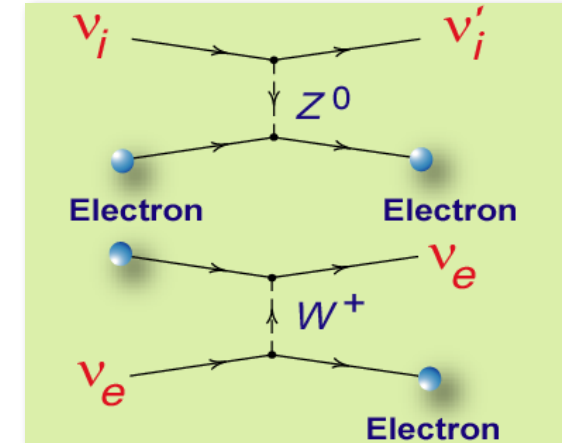
Why is this process not observed with  $\mu^-$  and  $\tau^-$  ?

## Neutral current (NC)



All flavours

## Elastic scattering (ES)

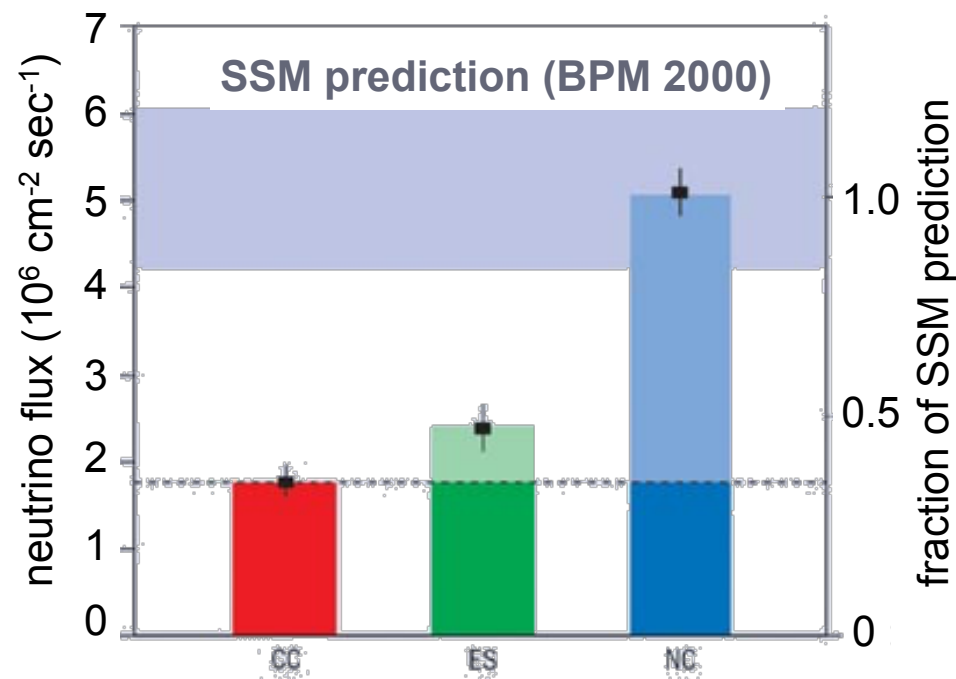


Mostly  $\nu_e$

- Scattering via neutral Z-Boson is flavour independent
- This reaction channel measures the entire neutrino flux
- NC detection enhanced by adding ~2 tons of salt (NaCl) to the heavy water (neutron capture on Cl nucleus, emitted gamma leads to detected signal)

# SNO results

- Results published in 2001 confirm flavour transformation hypothesis
- Neutral current reaction channel measures full neutrino flux expected in the Standard Solar Model
- Solar neutrino problem finally solved after 30 years!  
(Both Davis' experiment and Bahcall's calculations were right, after all ...)





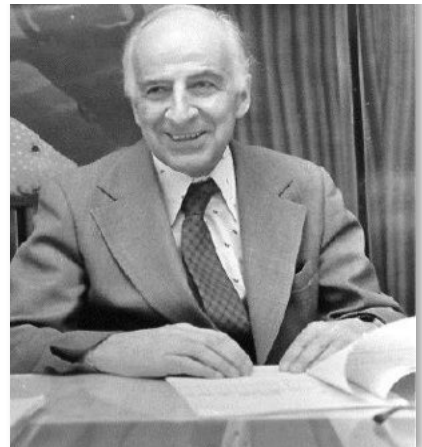
**Now that we're sure we see flavour oscillations:  
How can we explain them?**

# Neutrino oscillations

$\nu$ -oscillations are a **quantum mechanical interference phenomenon**

## 2-flavour mixing:

close analogy to CKM mixing of the left-handed quarks



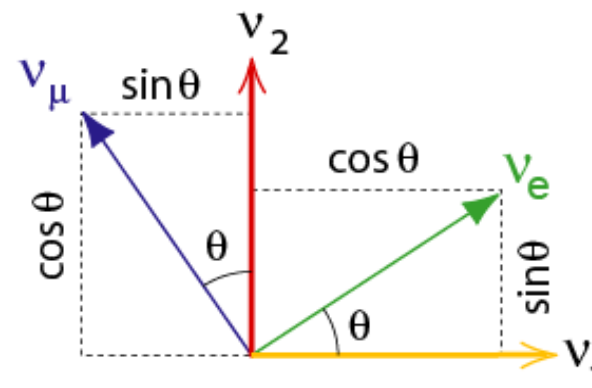
Bruno Pontecorvo:  
concept of  $\nu-\bar{\nu}$  oscillations

mass eigenstates

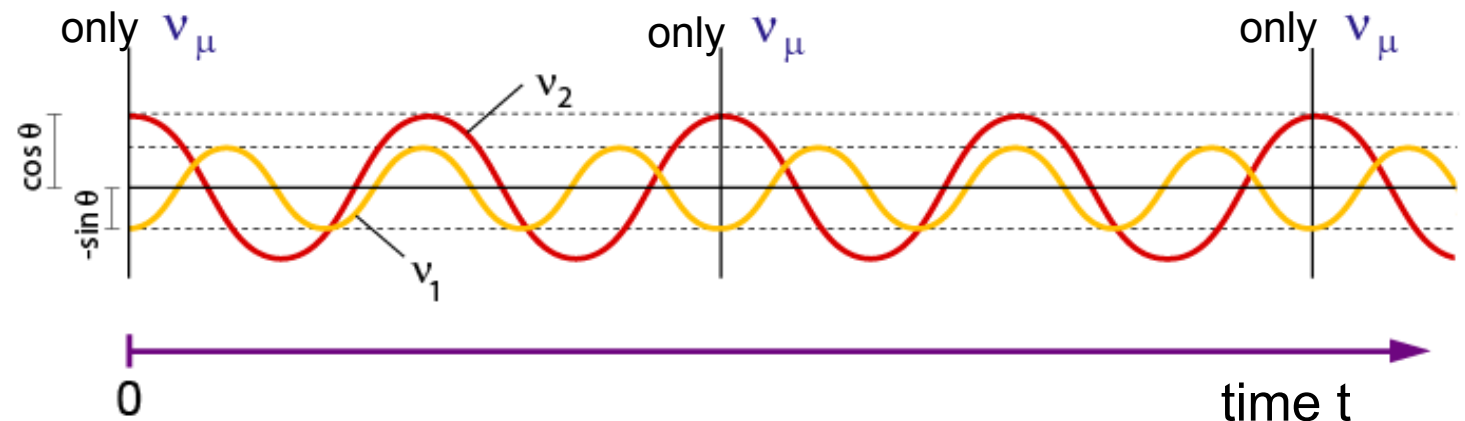


≠

flavour eigenstates



$$\begin{pmatrix} \nu_e \\ \nu_\mu \end{pmatrix} = \begin{pmatrix} \cos \theta & \sin \theta \\ -\sin \theta & \cos \theta \end{pmatrix} \begin{pmatrix} \nu_1 \\ \nu_2 \end{pmatrix}$$



# Neutrino oscillations

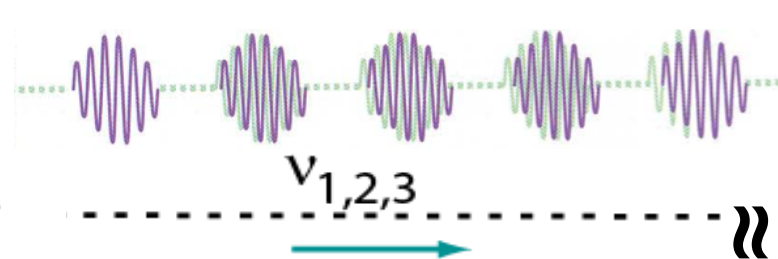
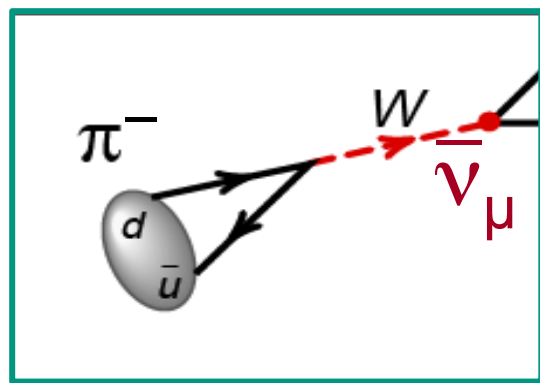
$\nu$ -oscillations result from different **propagation of mass eigenstates**

$\nu$  source

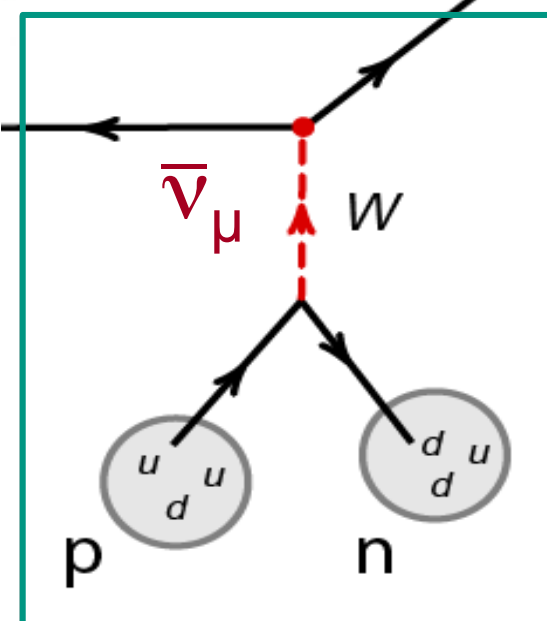
$\nu$  oscillations

$\nu$  detection

flavour state



flavour state  $\mu^+$



propagation of the  $\nu$ -mass eigenstates

$L = 10 \text{ m} \dots 10.000 \text{ km}$

known source:

- $\nu$  energies
- $\nu$  fluxes
- $\nu$  species

- detection efficiency
- energy resolution



# Neutrino oscillations – formalism

Probability  $P$  for the **oscillation of a  $\nu_\mu$  into a  $\nu_e$**  after time  $t$ :

$$P(\nu_\mu \rightarrow \nu_e) = |\cos\theta \cdot \sin\theta \cdot (1 - e^{i\Delta m^2 t / 2E_\nu})|^2$$

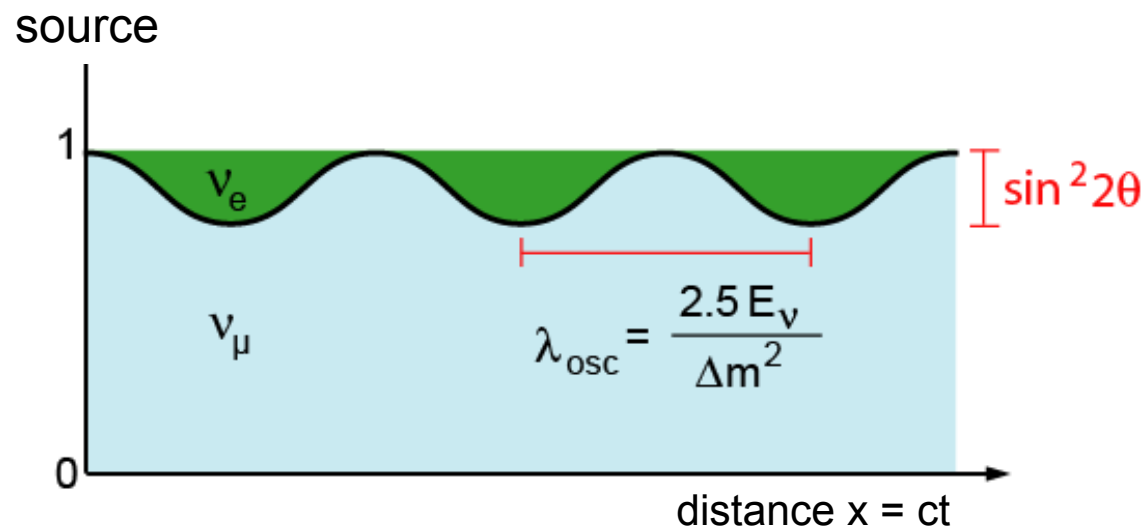
with  $P = |\langle \nu_e | \nu_\mu(t) \rangle|^2$

$$= \sin^2 2\theta \cdot \sin^2 (\Delta m^2 L_\nu / 4E_\nu)$$

with  $\Delta m^2 = |m_1^2 - m_2^2|$

amplitude      frequency

Oscillations only occur if at least one neutrino has a mass!

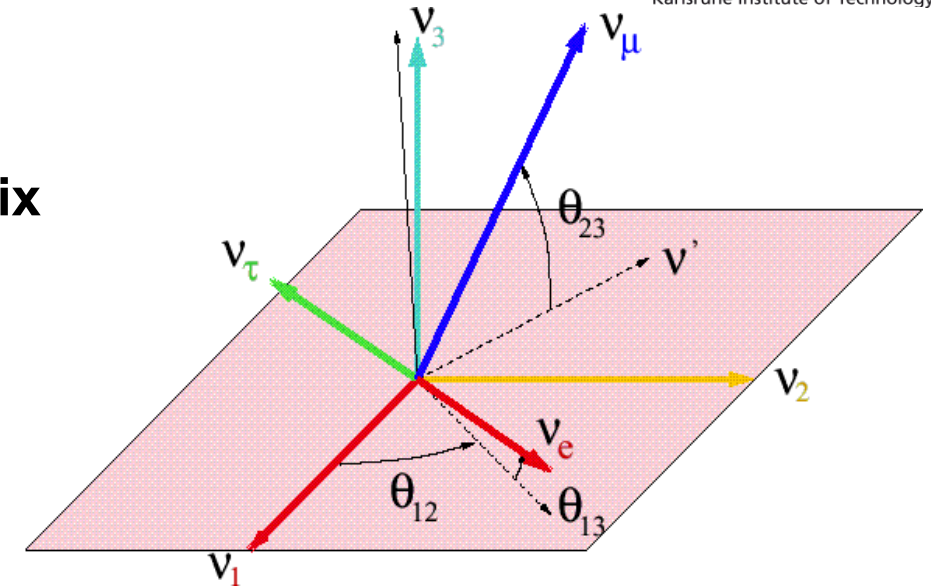


- Periodic decrease/increase of **primary** neutrino flavour state
- Oscillation length  $\lambda_{\text{osc}} \sim 2.5 E_\nu / \Delta m^2$
- Choose  $L$  and  $E$  such that you're sensitive to a given  $\theta$  and  $\Delta m^2$  and measure  $P(L/E)$

# The full three-flavour picture



3 x 3 unitary mixing matrix  
analogous to CKM:  
“Pontecorvo Maki  
Nakagawa Sakata”  
(PMNS)



- 3 mixing angles:  $\theta_{12}$ ,  $\theta_{23}$ ,  $\theta_{13}$ ,
- 1 CP-violating phase:  $\delta$
- two independent  $\Delta m^2$  scales:

$$\Delta m_{13}^2 = \Delta m_{12}^2 + \Delta m_{23}^2$$

$$U = \begin{pmatrix} 1 & 0 & 0 \\ 0 & \cos\theta_{23} & \sin\theta_{23} \\ 0 & -\sin\theta_{23} & \cos\theta_{23} \end{pmatrix} \cdot \begin{pmatrix} \cos\theta_{13} & 0 & \sin\theta_{13}e^{-i\delta} \\ 0 & 1 & 0 \\ -\sin\theta_{13}e^{-i\delta} & 0 & \cos\theta_{13} \end{pmatrix} \cdot \begin{pmatrix} \cos\theta_{12} & \sin\theta_{12} & 0 \\ -\sin\theta_{12} & \cos\theta_{12} & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

$\delta$ : CP-Phase

# Three-flavour neutrino mixing

## — present experimental values

$$U = \begin{pmatrix} 1 & 0 & 0 \\ 0 & \cos\theta_{23} & \sin\theta_{23} \\ 0 & -\sin\theta_{23} & \cos\theta_{23} \end{pmatrix} \cdot \begin{pmatrix} \cos\theta_{13} & 0 & \sin\theta_{13}e^{-i\delta} \\ 0 & 1 & 0 \\ -\sin\theta_{13}e^{-i\delta} & 0 & \cos\theta_{13} \end{pmatrix} \cdot \begin{pmatrix} \cos\theta_{12} & \sin\theta_{12} & 0 \\ -\sin\theta_{12} & \cos\theta_{12} & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

2. & 3. generation	1. & 3. generation	1. & 2. generation
$\Delta m_{23}^2 = 2.5 \times 10^{-3} \text{ eV}^2$	$\Delta m_{13}^2 = 2.5 \times 10^{-3} \text{ eV}^2$	$\Delta m_{12}^2 = 7.6 \times 10^{-5} \text{ eV}^2$
$\pm 6\%$	$\pm 6\%$	$\pm 7\%$
$\theta_{23} \approx 45^\circ$ ( <i>maximal?</i> )	$\theta_{13} \approx 8.5^\circ$ ( <i>small</i> )	$\theta_{23} \approx 34^\circ$ ( <i>large</i> )
$\pm 16\%$	$\pm 5\%$	$\pm 7\%$

3 $\sigma$  uncertainty,  
including ordering

3 $\sigma$  uncertainty,  
including ordering

→ Structure of leptonic mixing matrix very different from CKM matrix:

$$U_{\text{PMNS}} = \begin{pmatrix} 0.800 \rightarrow 0.844 & 0.515 \rightarrow 0.581 & 0.139 \rightarrow 0.155 \\ 0.229 \rightarrow 0.516 & 0.438 \rightarrow 0.699 & 0.614 \rightarrow 0.790 \\ 0.249 \rightarrow 0.528 & 0.462 \rightarrow 0.715 & 0.595 \rightarrow 0.776 \end{pmatrix}$$



# How do we know all this?

→ collected “world data” from many different experiments

2. & 3. generation	1. & 3. generation	1. & 2. generation
$\Delta m_{23}^2 = 2.5 \times 10^{-3} \text{ eV}^2$	$\Delta m_{13}^2 = 2.5 \times 10^{-3} \text{ eV}^2$	$\Delta m_{12}^2 = 7.6 \times 10^{-5} \text{ eV}^2$
$\theta_{23} \approx 45^\circ$ (maximal?)	$\theta_{13} \approx 8.5^\circ$ (small)	$\theta_{23} \approx 34^\circ$ (large)

atmospheric  
& long-baseline  
accelerator exp.

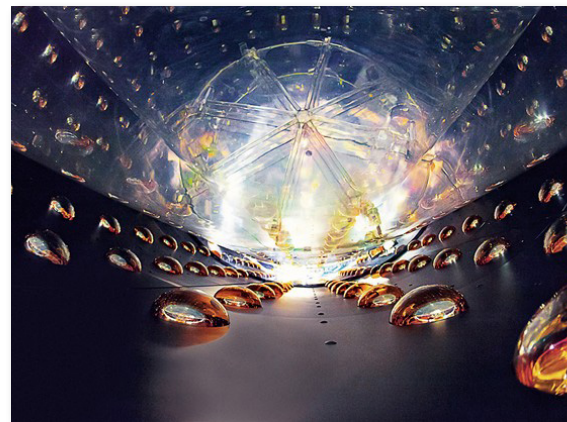
GeV,  $\nu_\mu$  ( $\bar{\nu}_\mu$ )



T2K (Tokai-Kamioka)

reactor &  
long-baseline  
accelerator exp.

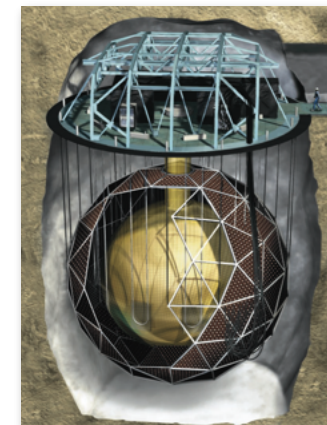
MeV,  $\bar{\nu}_e$   
GeV,  $\nu_\mu$  ( $\bar{\nu}_\mu$ )



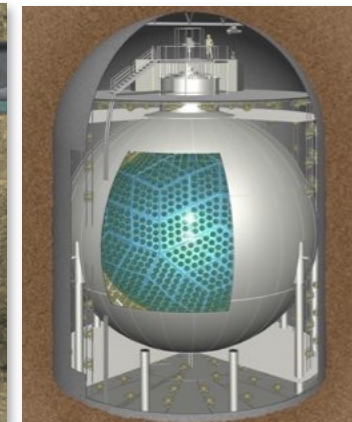
Daya Bay

solar  
& reactor exp.

MeV,  $\nu_e$  ( $\bar{\nu}_e$ )



SNO

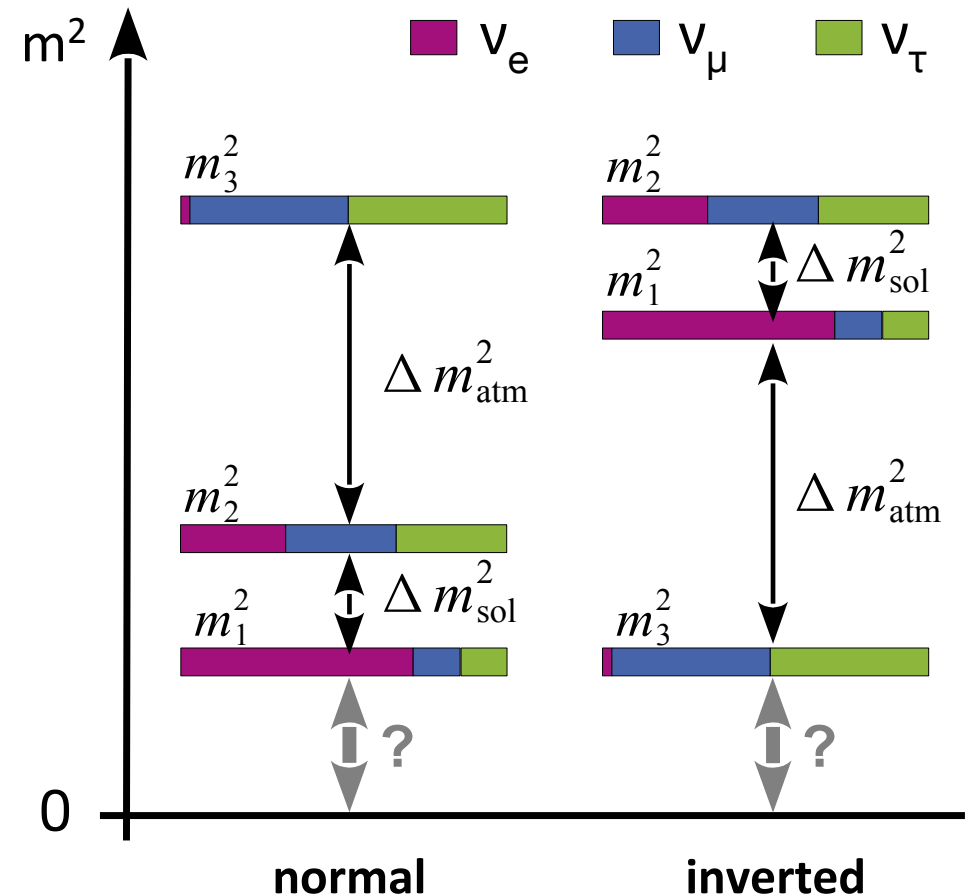


KamLAND

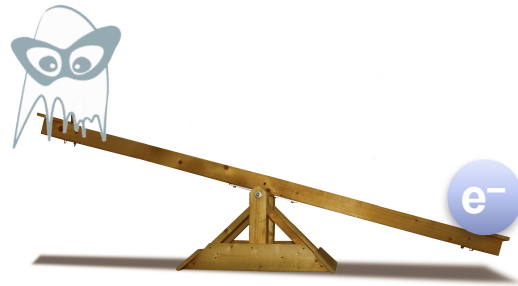
# The big picture: What have we learned from oscillation data?

- Large neutrino mixing and tiny neutrino masses  $m(\nu_i) \neq 0$  established
- Evidence for non-zero  $\theta_{13}$
- Hints for non-maximal  $\theta_{23} \neq \pi/4$
- Expectation of CP-violating phase  $\delta$
- Absolute mass scale cannot be determined from oscillations
- Expect  $m_\nu > 10$  meV for normal ordering,  $m_\nu > 50$  meV for inverted
- Majorana vs Dirac nature of neutrinos?

**New!  
BSM physics!**



## IV. How can we measure neutrino masses?



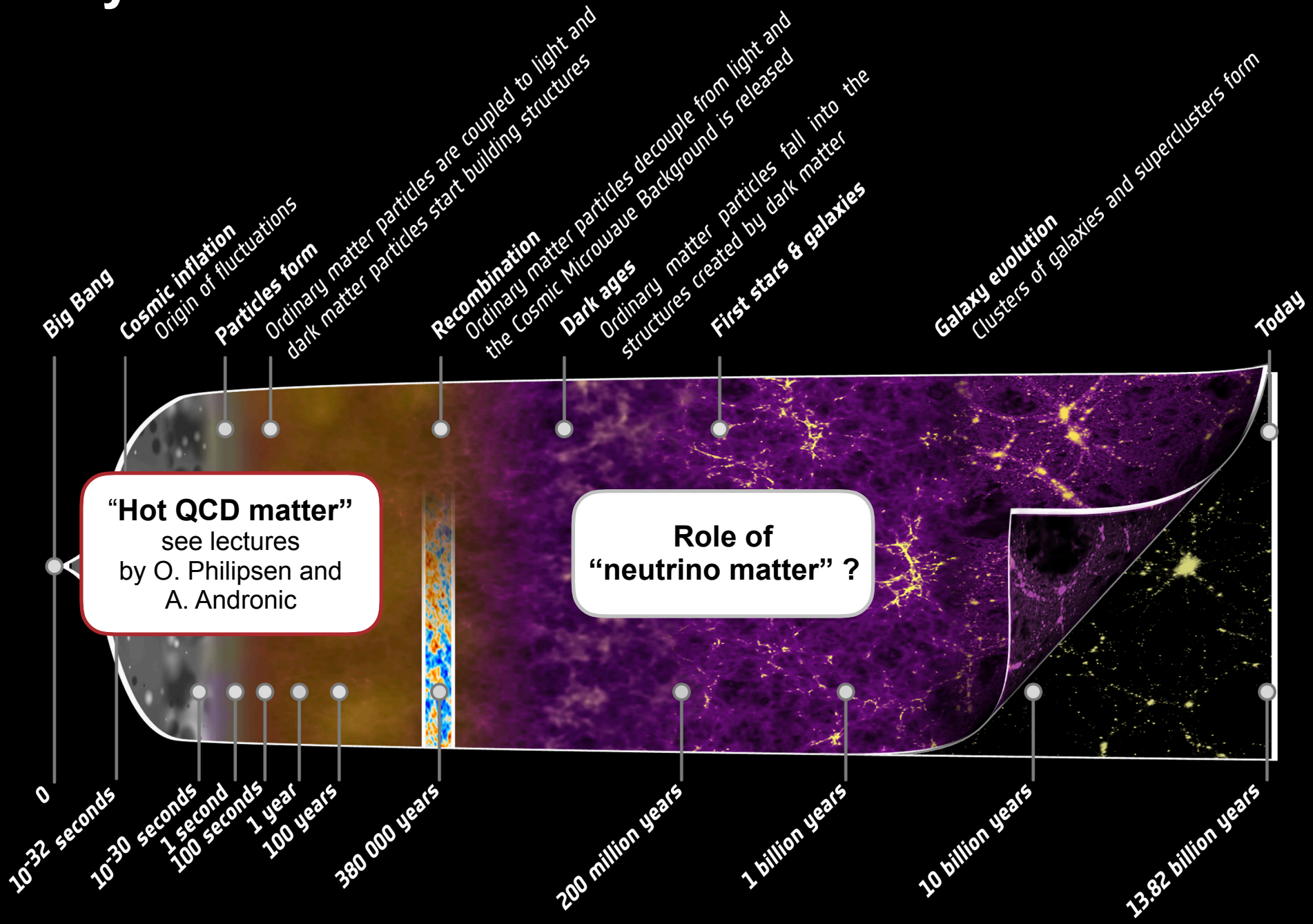
Indirect (model-dependent) probes:

- Observational cosmology
- Search for  $0\nu\beta\beta$

Direct (model-independent) probes:

- Kinematics of weak decays  
( ${}^3\text{H}$   $\beta$ -decay,  ${}^{163}\text{Ho}$  EC)

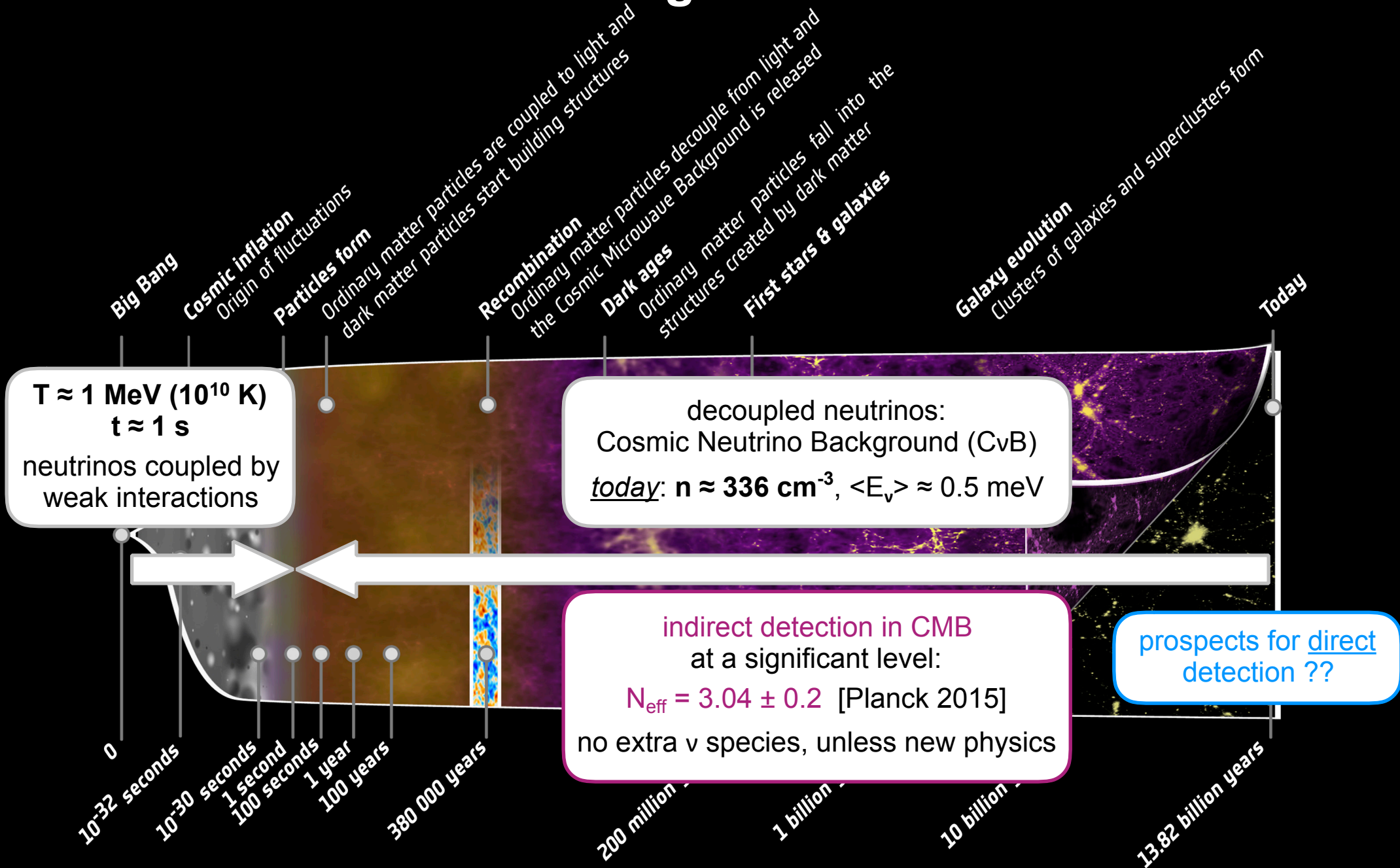
# History of the Universe



[ESA, Planck]



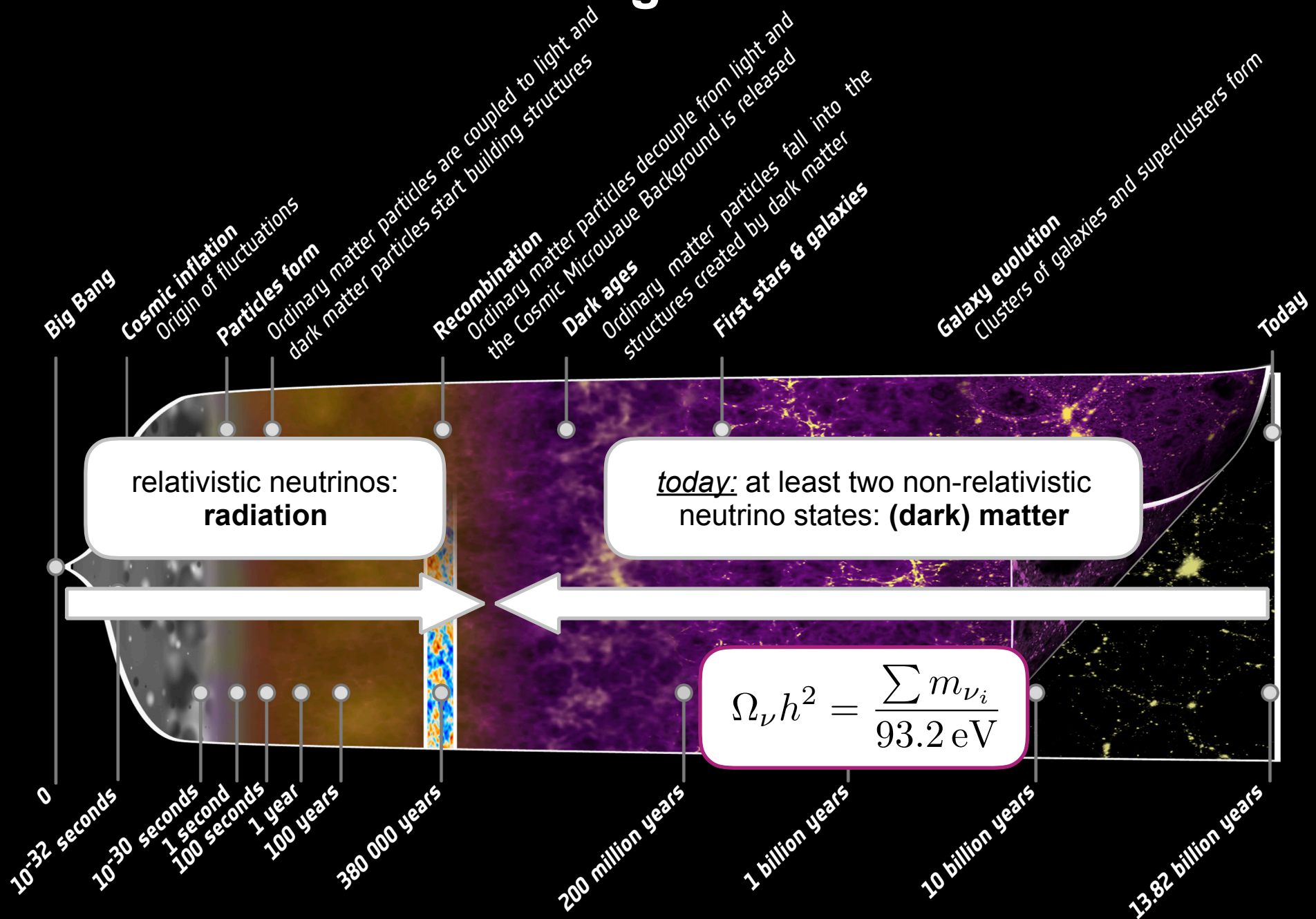
# The Cosmic Neutrino Background



[ESA, Planck]

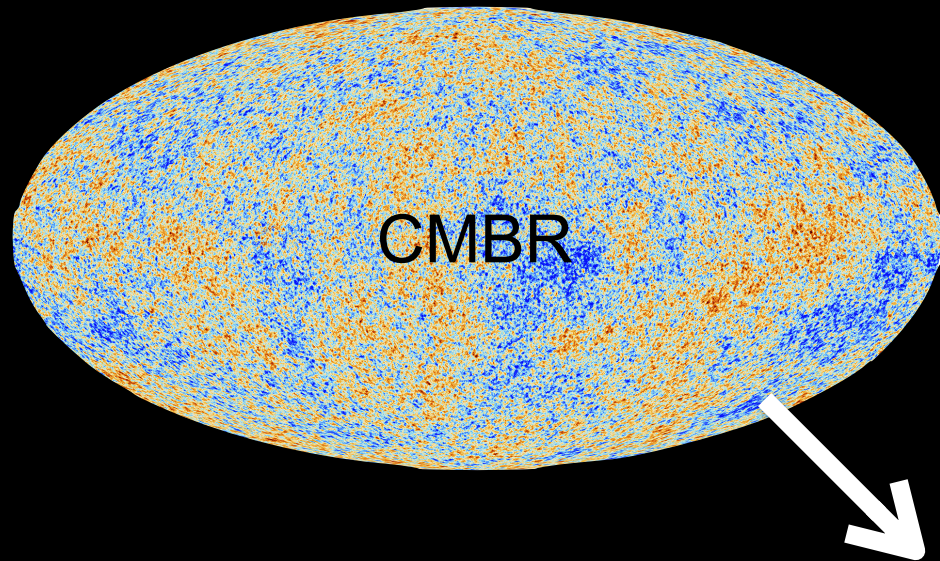


# The Cosmic Neutrino Background

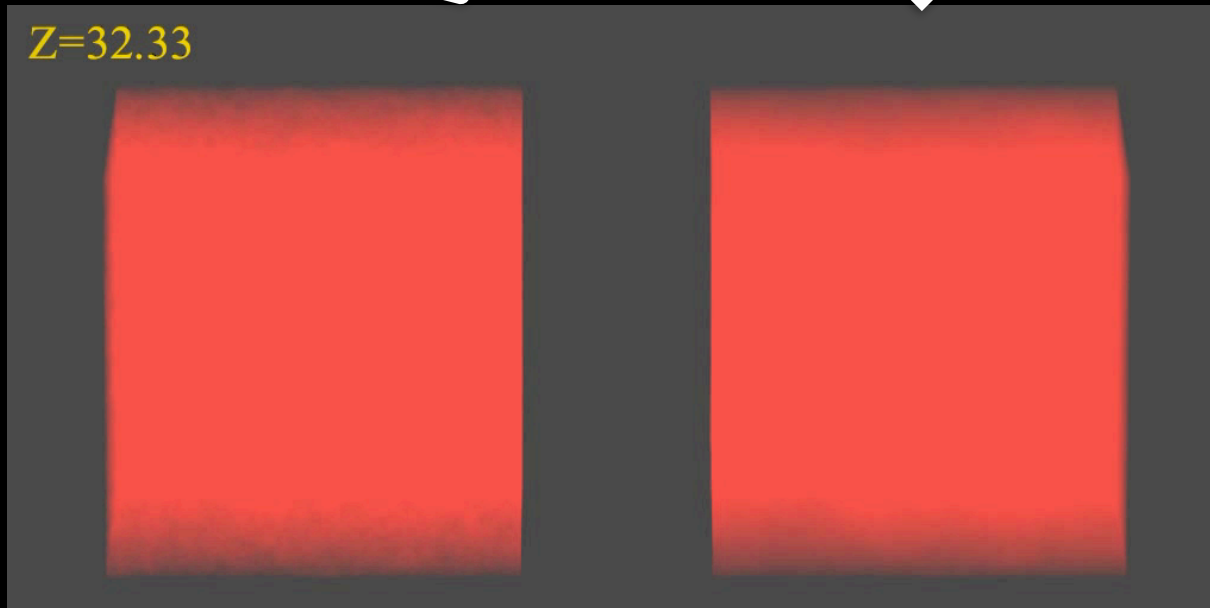
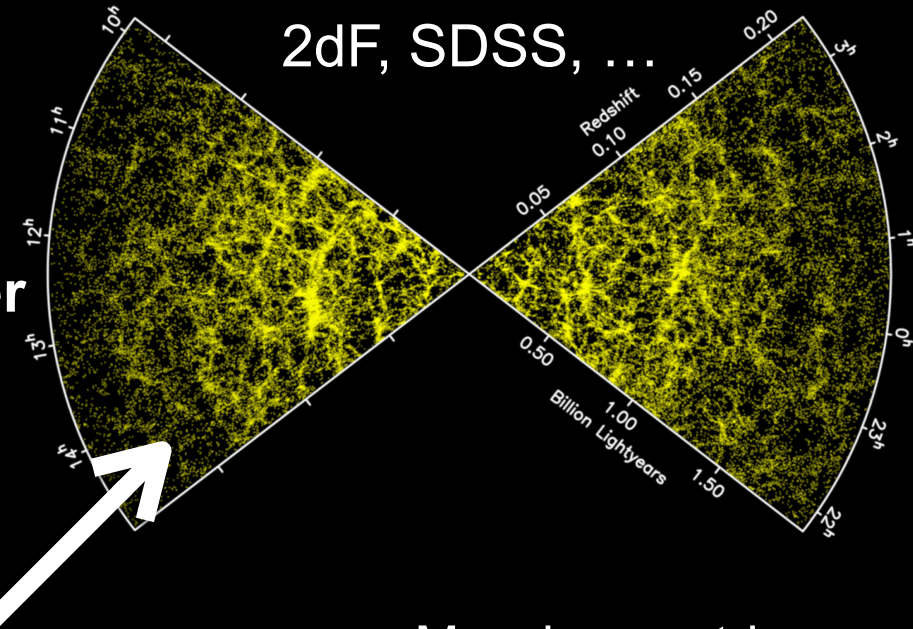


[ESA, Planck]

# Neutrino mass: adding information



cold dark matter  
+  
relic neutrino  
density:  
 $336 \nu / \text{cm}^3$



$m_\nu = 0 \text{ eV}$

$2 \text{ eV}$

[T. Haugbølle, Univ. of Aarhus]

- Massive neutrinos wash out structure at small scales

- Status 2017:

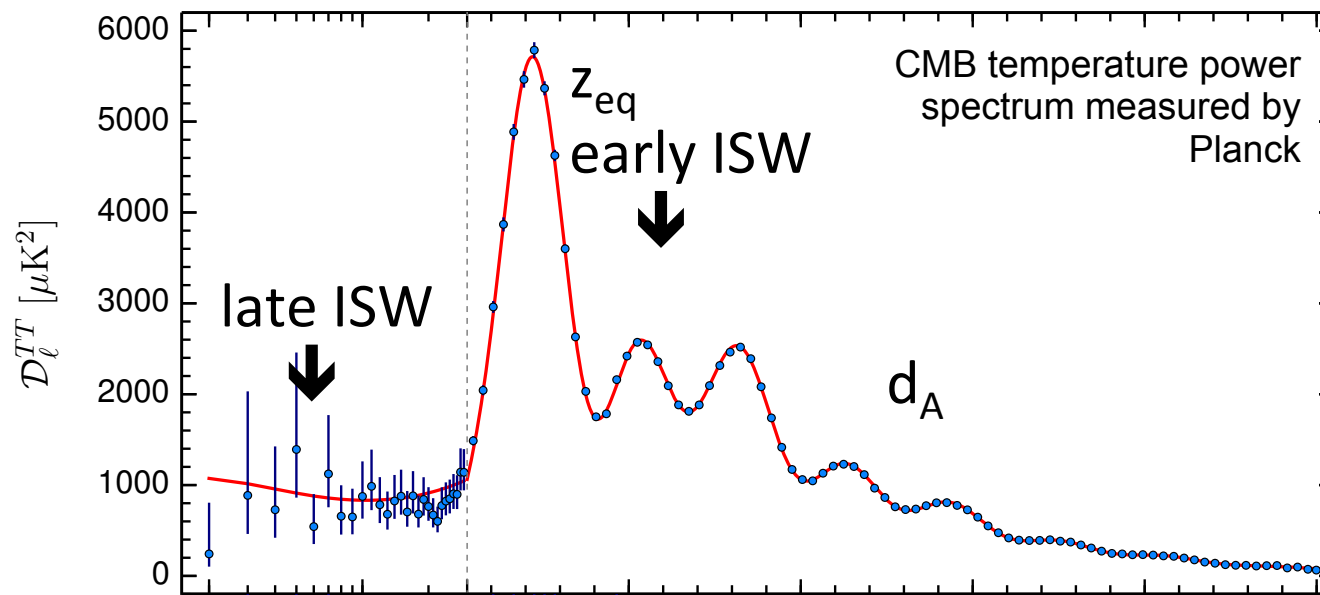
$$\sum m_\nu \lesssim 0.13 \text{ eV}$$

using CMB + LSS  
+ BAO

- Caveat:  
degeneracies

# Neutrino mass from cosmology

- Current **observational cosmology** offers a wealth of precision data which can be combined to learn about neutrino masses
- Requires interpretation in the framework of the **Standard Model ( $\Lambda$ CDM)** of Cosmology
- CMB measurements only (pre-Planck):
  - probe neutrino mass mainly via Integrated Sachs-Wolfe effect (modified grav. potential seen by photons)
  - neutrinos contribute to radiation density at  $z_{\text{eq}}$  and to non-rel. matter density today



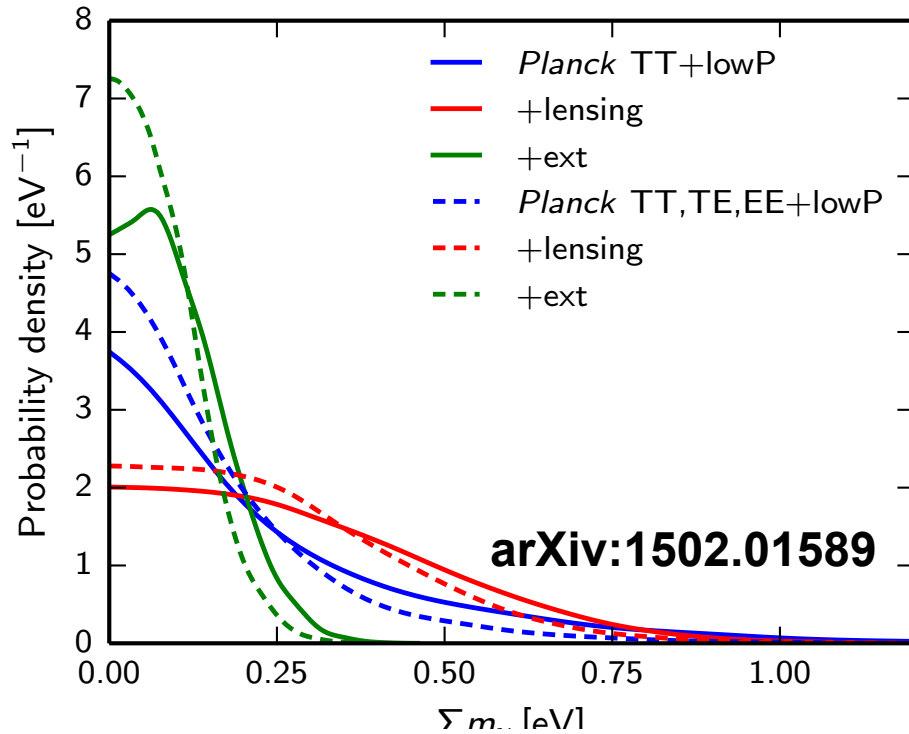
- Post-Planck:  
Weak lensing of CMB gives additional information on  $\Sigma m_\nu$
- Status 2017:

$$\sum m_\nu \lesssim 0.6 \text{ eV}$$

using only CMB temperature & polarization data



# Neutrino mass from cosmology: other probes



$$\begin{aligned} \Sigma m_\nu &< 0.72 \text{ eV} && Planck \text{ TT+lowP}; \\ \Sigma m_\nu &< 0.21 \text{ eV} && Planck \text{ TT+lowP+BAO}; \\ \Sigma m_\nu &< 0.49 \text{ eV} && Planck \text{ TT, TE, EE+lowP}; \\ \Sigma m_\nu &< 0.17 \text{ eV} && Planck \text{ TT, TE, EE+lowP+BAO}. \end{aligned}$$

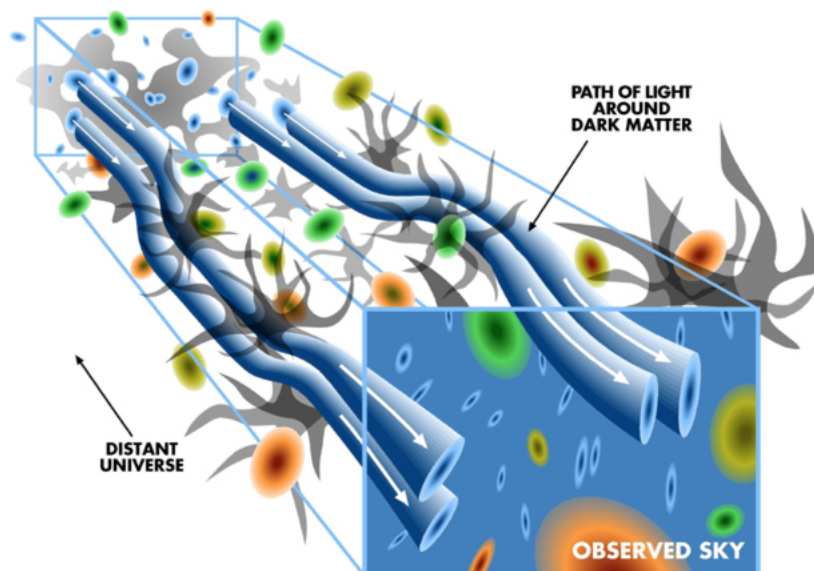
**6 + 2 parameter  $\Lambda$ CDM model**

reach:  
20-50 meV

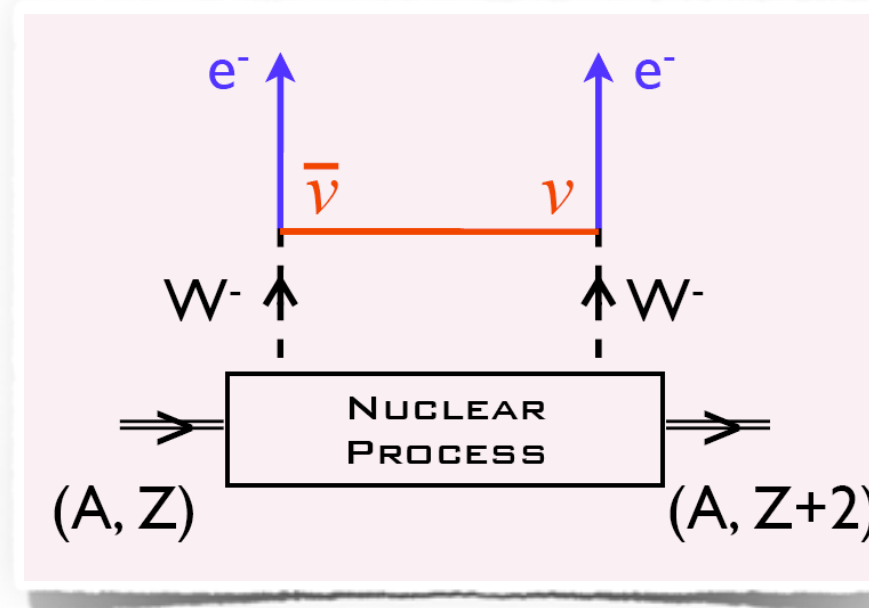
**Future EUCLID mission (ESA):  
grav. lensing & galactic power spectrum**

“... *Euclid* will very likely provide a **positive detection** of neutrino mass ..., the exact nature of the neutrino mass spectrum remains out of its reach ...”

[Hamann, Hannestad, & Wong, JCAP 11 (2012) 52]



## Search for neutrinoless double beta decay



Are neutrinos Majorana fermions ( $\nu = \bar{\nu}$ ) ?

Is lepton number violated ( $\Delta L = 2$ ) ?



# Double $\beta$ -decay: $2\nu\beta\beta$ & $0\nu\beta\beta$ modes

## Double beta decay with neutrino emission ( $2\nu\beta\beta$ ):

2nd-order weak interaction process

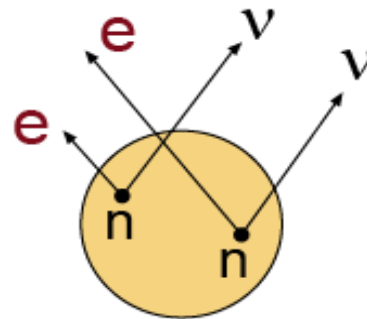
↪ extremely small transition rates & long half-lives:  $T_{1/2} \sim 10^{19} - 10^{21}$  years

↪ energy  $E_0$  shared by 4 leptons, observed in 12 isotopes so far

$2\nu\beta\beta$

$$(Z, A) \rightarrow (Z + 2, A)$$

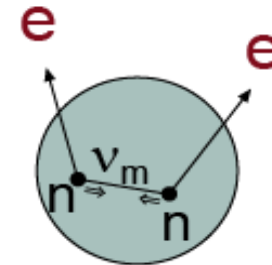
$$+ e_1^- + e_2^- + \bar{\nu}_{e,1} + \bar{\nu}_{e,2}$$



$0\nu\beta\beta$

$$(Z, A) \rightarrow (Z + 2, A)$$

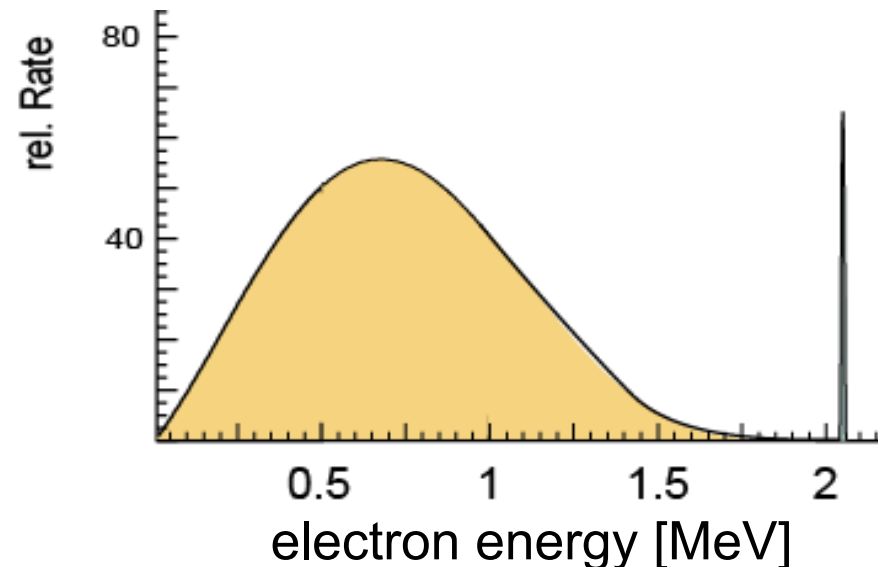
$$+ e_1^- + e_2^-$$



first description:



M. Goeppert-Mayer (1935)



first description:



G. Racah (1937)



E. Majorana

# Double beta decay: candidate nuclides

35 “energetically” suitable even-even nuclides for double beta decay

Decay rates  $\sim Q^5$  ( $0\nu\beta\beta$ )  $\rightarrow$  find suitable isotope for experiment:

**11 nuclei for  $2\nu\beta^-\beta^-$  at  $Q > 2$  MeV:**

$\beta^-\beta^-$ decay	Q [MeV]	nat. [%]
$^{48}\text{Ca} \rightarrow ^{48}\text{Ti}$	4,274	0,187
$^{76}\text{Ge} \rightarrow ^{76}\text{Se}$	2,039	7,8
$^{82}\text{Se} \rightarrow ^{82}\text{Kr}$	2,995	9,2
$^{96}\text{Zr} \rightarrow ^{96}\text{Mo}$	3,348	2,8
$^{100}\text{Mo} \rightarrow ^{100}\text{Ru}$	3,034	9,6
$^{110}\text{Pd} \rightarrow ^{110}\text{Cd}$	2,004	11,8
$^{116}\text{Cd} \rightarrow ^{116}\text{Sn}$	2,809	7,5
$^{124}\text{Sn} \rightarrow ^{124}\text{Te}$	2,288	5,64
$^{130}\text{Te} \rightarrow ^{130}\text{Xe}$	2,527	34,5
$^{136}\text{Xe} \rightarrow ^{136}\text{Ba}$	2,458	8,9
$^{150}\text{Nd} \rightarrow ^{150}\text{Sm}$	3,368	5,6

**6 nuclei for  $2\nu\beta^+\beta^+/\text{EC}$  at lower Q:**

$\beta^+\beta^+$ decay	Q [MeV]	nat. [%]
$^{78}\text{Kr} \rightarrow ^{78}\text{Se}$	<b>0,838</b>	0,35
$^{96}\text{Ru} \rightarrow ^{96}\text{Mo}$	<b>0,676</b>	5,5
$^{106}\text{Cd} \rightarrow ^{106}\text{Pd}$	<b>0,738</b>	1,25
$^{124}\text{Xe} \rightarrow ^{124}\text{Te}$	<b>0,822</b>	0,10
$^{130}\text{Ba} \rightarrow ^{130}\text{Xe}$	<b>0,534</b>	0,11
$^{136}\text{Ce} \rightarrow ^{136}\text{Ba}$	<b>0,362</b>	0,19

Coulomb barrier reduces Q  
 $\rightarrow$  even longer expected  $T_{1/2}$

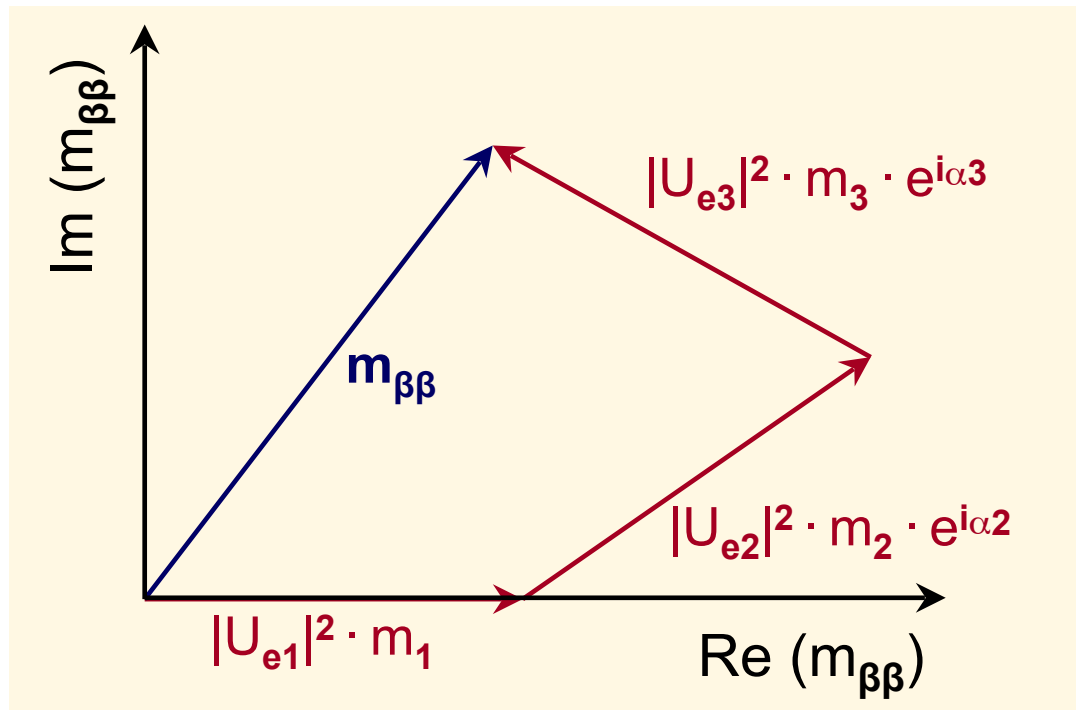
$$T_{1/2} (\beta^+\beta^+) \sim 10^{26} \text{ a}$$

# Double $\beta$ -decay & Majorana mass $m_{\beta\beta}$

**Effective Majorana mass  $m_{\beta\beta}$**  is not identical with  $m(\nu_e)$  from  $\beta$ -decay  
**Coherent** sum over three  $\nu$  mass eigenstates  $m_1, m_2, m_3$

$$\langle m_{\beta\beta} \rangle = \left| \sum_{i=1}^3 |U_{e,i}|^2 m_i \cdot e^{i\alpha_i} \right|$$

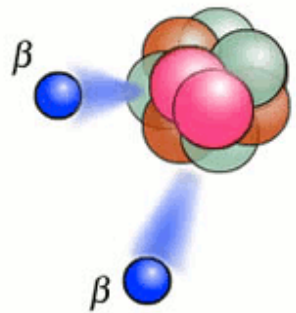
2 independent Majorana CP-phases  $\alpha_i$   
 $\Rightarrow$  mutual cancellations are possible  
 if  $\alpha_i \neq n \cdot \pi \rightarrow$  CP violation



virtual particles interfere!

# Double $\beta$ -decay & Majorana mass $m_{\beta\beta}$

Determination of **effective Majorana mass  $m_{\beta\beta}$**  from  $0\nu\beta\beta$  half-life  $T_{1/2}$



$$\langle m_{\beta\beta} \rangle^2 = \left( T_{1/2}^{0\nu\beta\beta} \cdot G^{0\nu\beta\beta}(E_0, Z) \cdot \left| M_{GT}^{0\nu\beta\beta} - \frac{g_V^2}{g_A^2} M_F^{0\nu\beta\beta} \right|^2 \right)^{-1}$$

experimental value of  
 $0\nu\beta\beta$  **half-life**

**phase space factor**

**nuclear matrix element**

$M_{GT}$ : Gamov-Teller

$M_F$ : Fermi

- **experimental observable:  $T_{1/2}$**

↳  $0\nu\beta\beta$  event number depending on measuring time, number of target nuclei, experimental efficiency, background

- **weak interaction (phase space factor):  $G^{0\nu\beta\beta}$**

↳ determined by  $\beta\beta$ -endpoint energy; strong dependence  $\sim Q^5$

- **nuclear physics (matrix elements):  $M^{0\nu\beta\beta}$**

↳ shell model calculations, large uncertainties **O(100%)**

# Sensitivity drivers

$$T_{1/2}^{0\nu}(\text{FOM}) \propto a \cdot \epsilon \cdot \sqrt{\frac{M \cdot t}{b \cdot \Delta E}}$$

## Requirements:

- Large isotopical abundance (***a***)
- High efficiency (***ε***)
- Large Mass (***M***)
- Long counting time (***t***)
- Low background (***b***)
- Good energy resolution ( **$\Delta E$** )

→ Many suitable combinations for isotope + detector technology

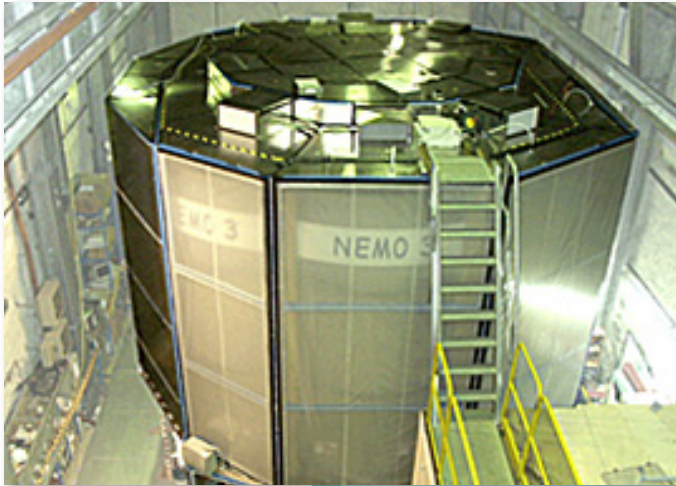
→ If ROI is background free:  
linear scaling with ***M*** and ***t***!

$$T_{1/2}^{0\nu}(\text{FOM}) \propto a \cdot \epsilon \cdot M \cdot t$$



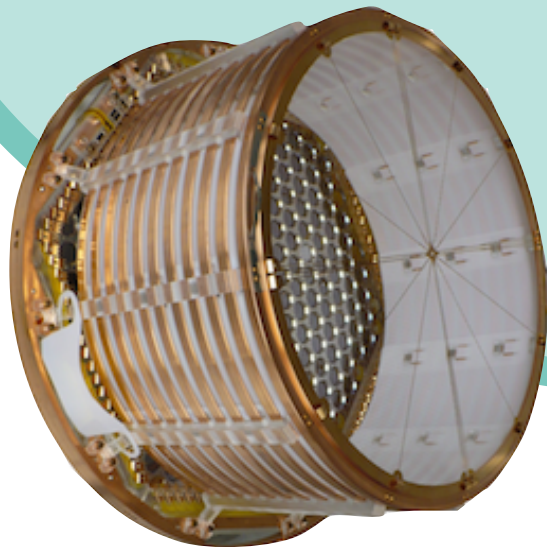
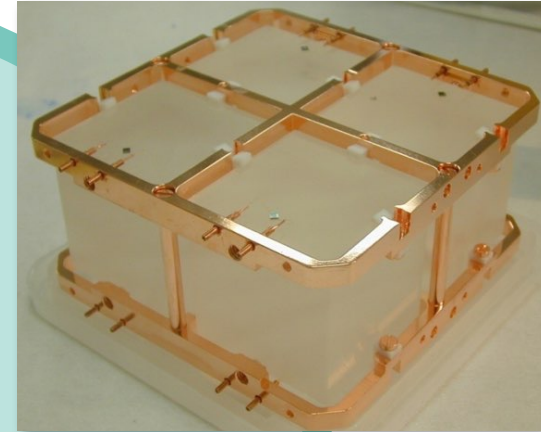
# Experimental techniques

Tracker-calorimeter

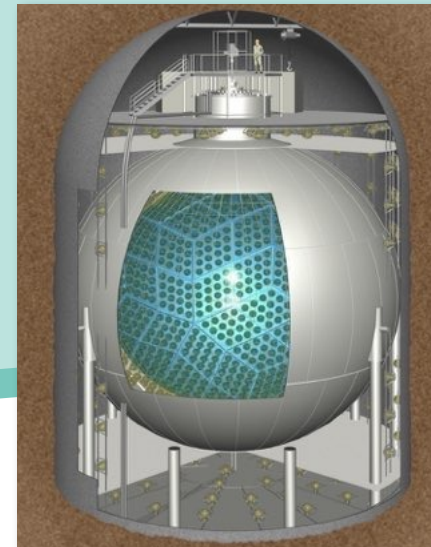


Ge diodes

Cryo-bolometer



Liquid noble element TPC



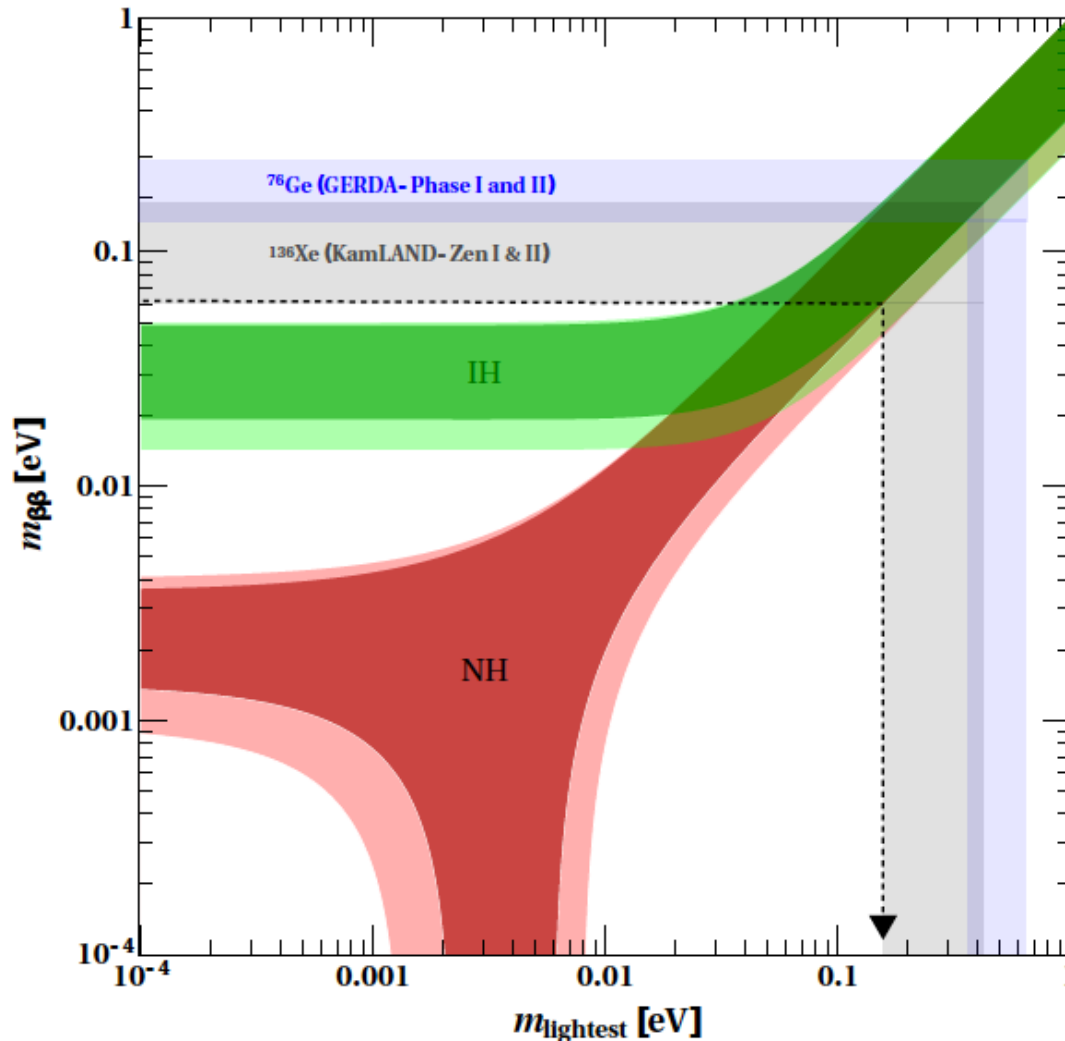
Liquid scintillator

# Many experimental ideas ...

Isotope	past generation	future generation	type
$^{76}\text{Ge}$	GERDA / MAJORANA	LEGEND	semiconductor detectors
$^{82}\text{Se}$	NEMO-3	SuperNEMO	tracking calorimeters
$^{130}\text{Te}$	CUORE	CUPID	bolometers/scintillators (diff isot considered for CUPID)
$^{130}\text{Te}$		SNO+	liquid scintillator
$^{136}\text{Xe}$	KamLAND-Zen	KamLAND2-Zen	liquid scintillator
$^{136}\text{Xe}$	EXO-200	nEXO	liquid TPC
$^{136}\text{Xe}$		NEXT/PANDA-X III	gas TPC

+ further projects in R&D phase

# Current constraints



Most stringent bounds now approaching inverted hierarchy

Next generation has good discovery potential, even for normal hierarchy



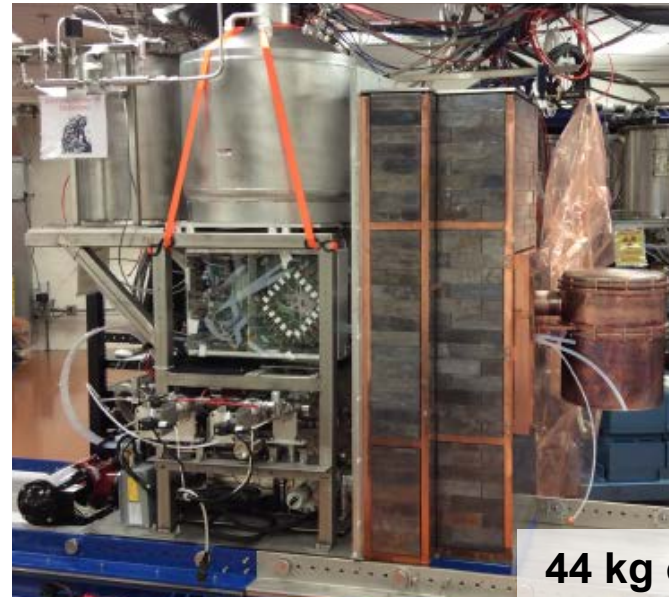
# Germanium diodes: MAJORANA and GERDA

## MAJORANA

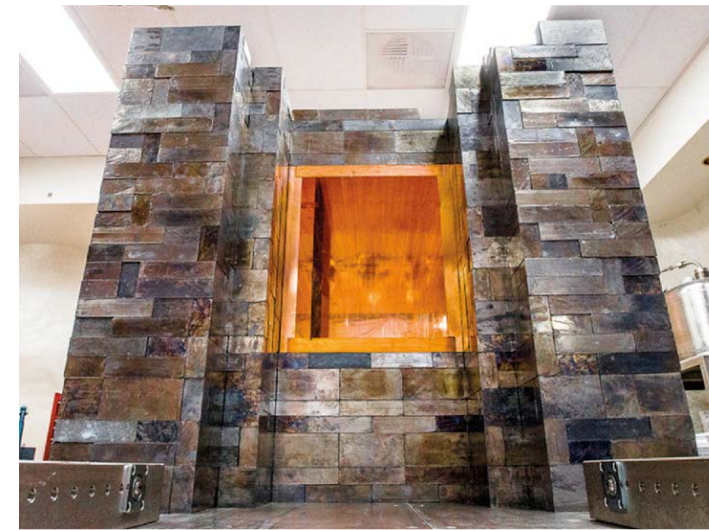


*Conventional design:*

*Vacuum cryostats in a passive graded shield with ultra-clean materials*



44 kg of Ge crystals (88%  $^{76}\text{Ge}$ ),  $\Delta E \sim 0.1\%$



SURF (South Dakota, USA)

## GERDA



*Novel design:*

*Direct immersion in active LAr shield*



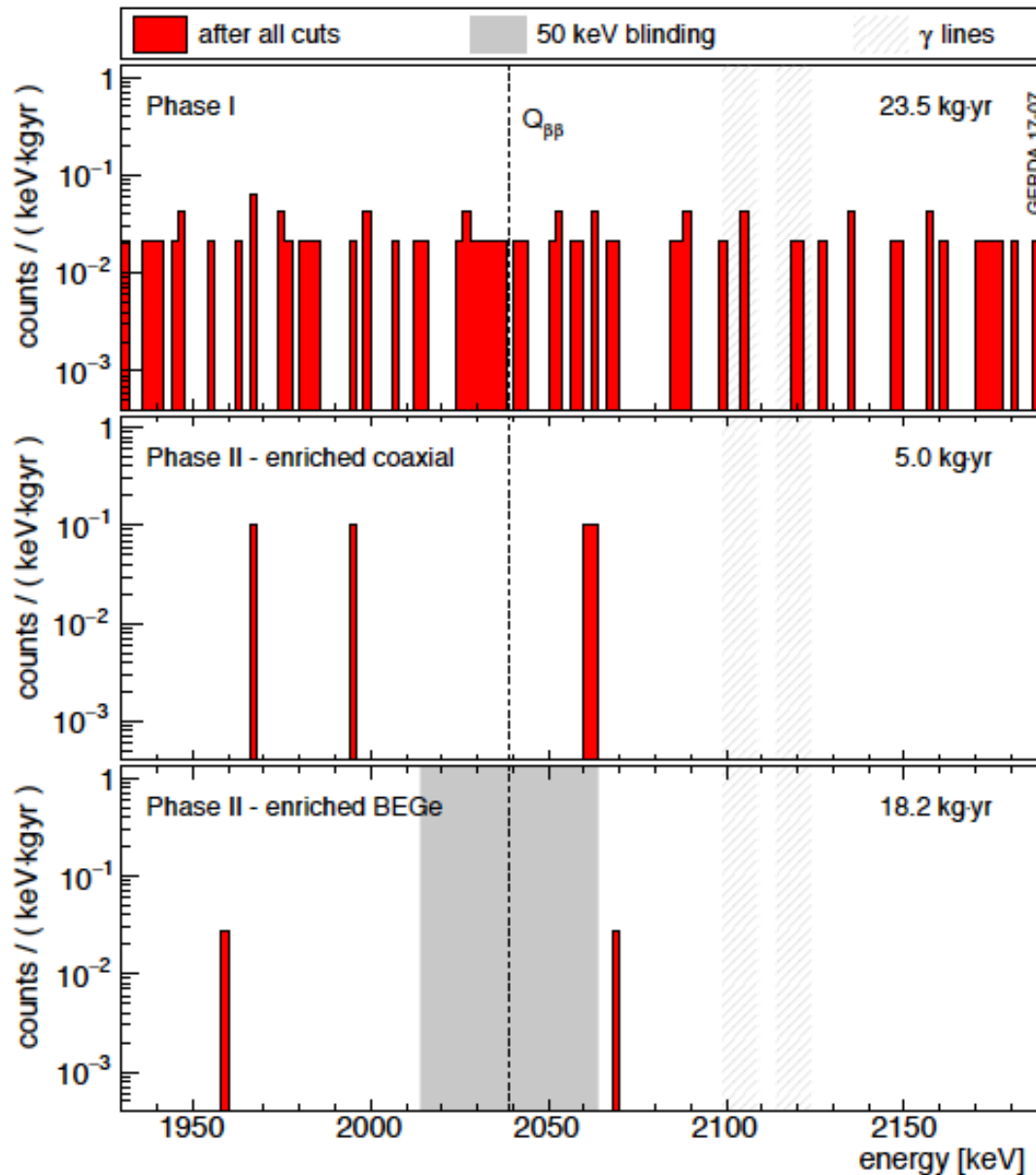
36 kg of Ge crystals (87%  $^{76}\text{Ge}$ ),  $\Delta E \sim 0.2\%$



Gran Sasso (Italy)

Alan Poon (LBNL), Erice 2017

# Germanium diodes: results

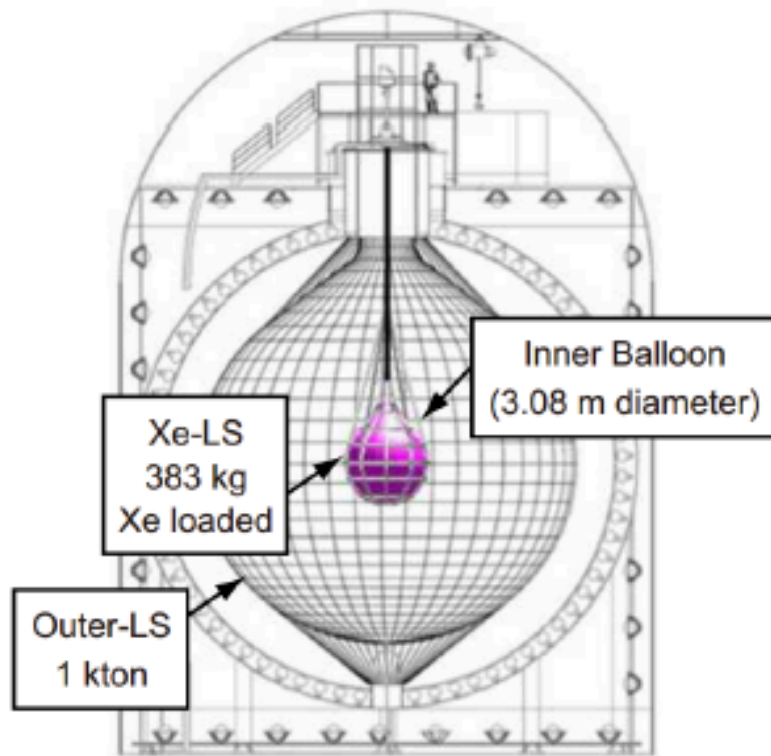


- Combined analysis **GERDA phase I + II**
- “Background-free” running in phase II
  - Two counts after unblinding
  - No count at  $Q_{\beta\beta}$
- $T_{1/2} > 8.0 \times 10^{25}$  yr (90% CL)
- Next-generation project: LEGEND  
 “Large Enriched Germanium Experiment for Neutrinoless  $\beta\beta$  Decay”
- Staged approach, starting with ~200 kg in existing GERDA cryostat
- Final goal: 1000 kg-scale detector for sensitivity  $>10^{27}$  yr
- Background improvement required: x30 (x5 for LEGEND-200)



# Kamland-Zen at Kamioka (Japan)

Isotope:	$^{136}\text{Xe}$ ( $Q_{\beta\beta}=2458$ keV)
Resolution:	240-270 keV FWHM
Mass:	350 kg
Technology:	Xe-loaded liquid scintillator
Status:	completed/upgrading

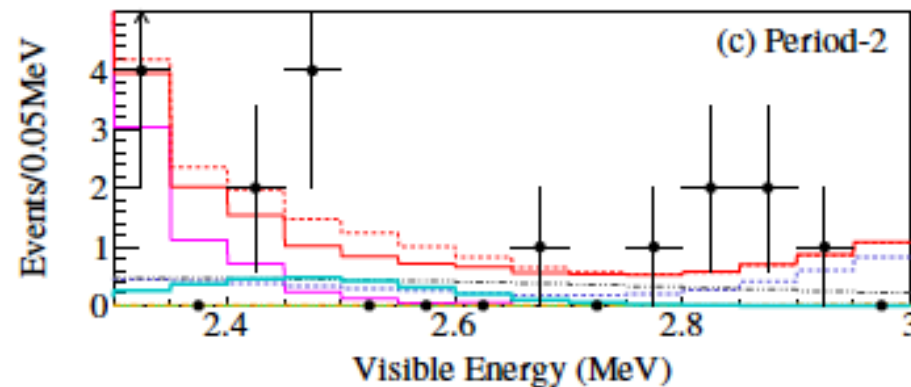
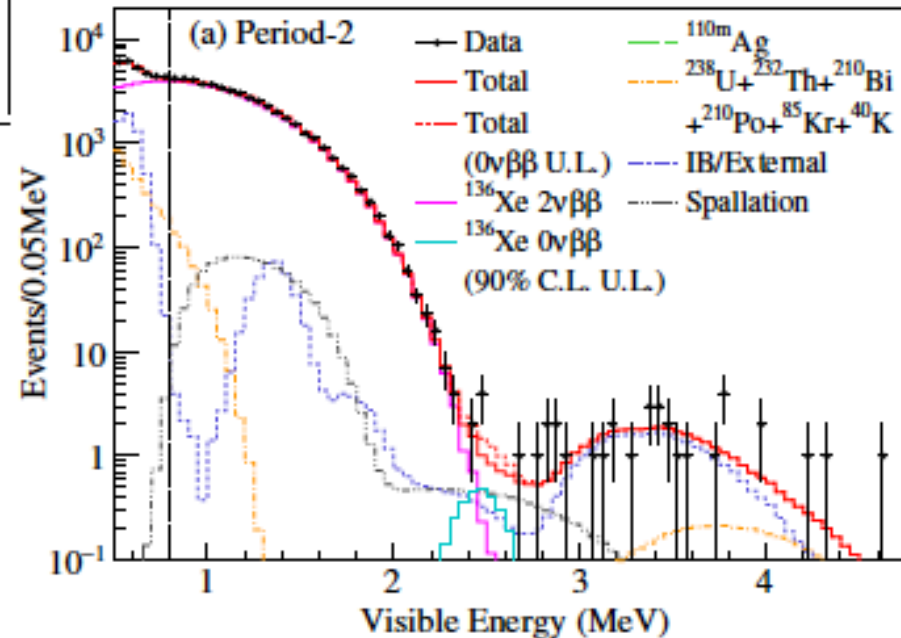


[Phys.Rev.Lett. 117 (2016) no.10, 109903]

Latest result with 504 kg·yr + old data:

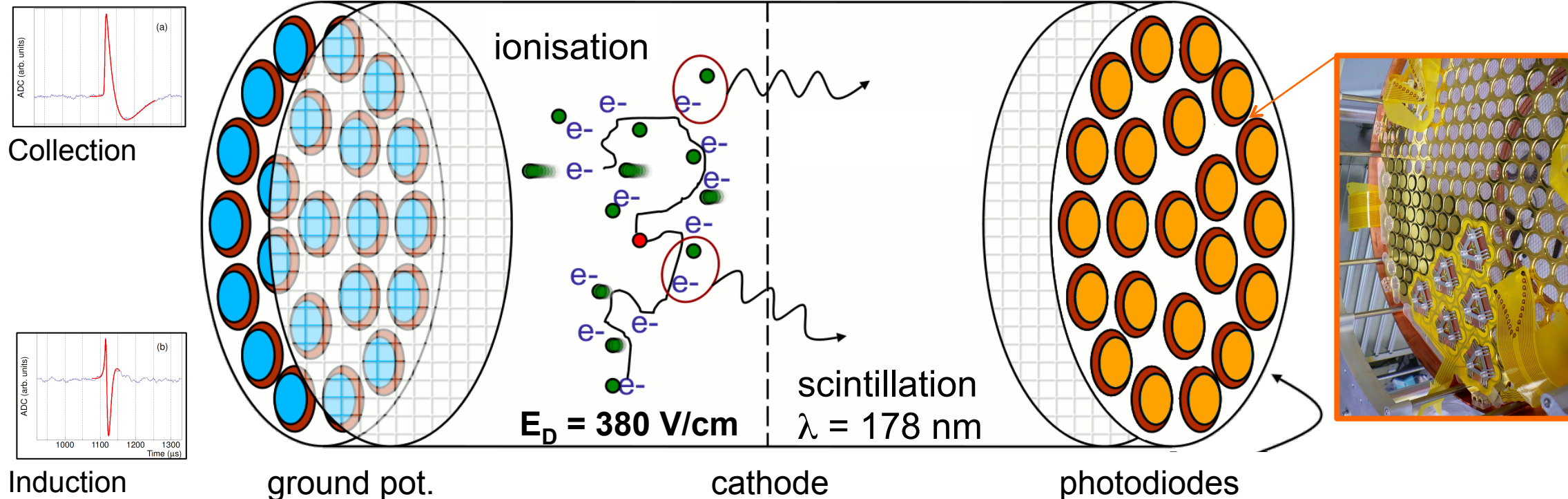
$$T_{1/2}^{0\nu} > 10.7 \cdot 10^{25} \text{ yr (90\% C.L.)}$$

$$T_{1/2}^{0\nu} > 5.6 \cdot 10^{25} \text{ yr (sensitivity)}$$



# Liquid Xenon TPC: EXO-200

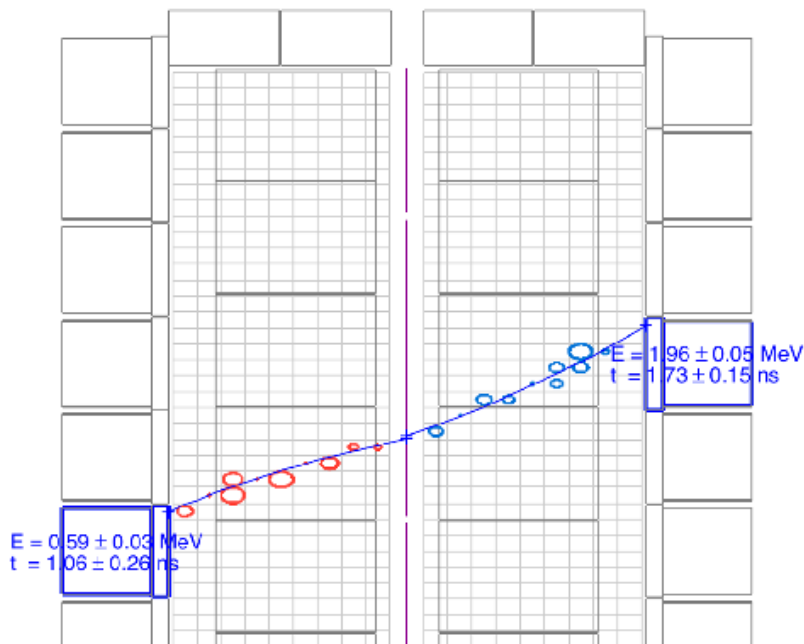
- Enriched Xenon Observatory at WIPP/New Mex., running **~175 kg of LXe** (80.6%  $^{136}\text{Xe}$ )
- More than a calorimeter: spatial resolution (x,y,z) and PID allows discrimination of multi-site (bg-like) vs. single-site ( $0\nu\beta\beta$ -like) events
- Anticorrelation of charge and light signals (compare DM detectors), tags  $\alpha$  events
- Now preparing **nEXO**: 5-ton monolithic detector ( $\sim 1$  t fiducial), 1.3 m electron drift length,  $\sim 4$  m<sup>2</sup> of SiPM photosensors, option of  $^{136}\text{Ba}$  tagging



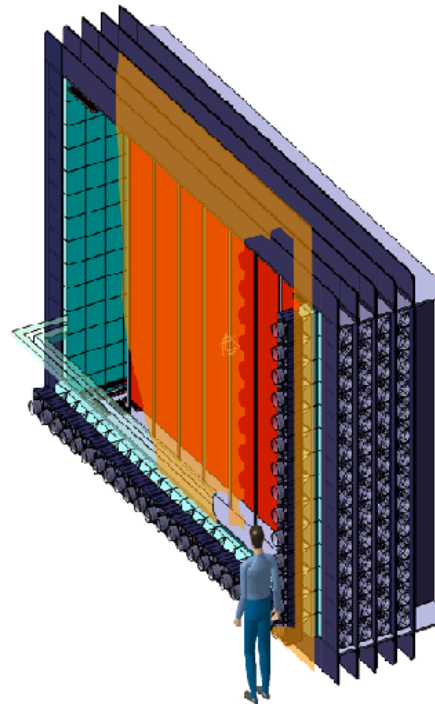
# Tracking-Calorimeter: SuperNEMO

- Successor of NEMO-3 at Laboratoire Souterrain de Modane (LSM)
- Baseline isotope:  $^{82}\text{Se}$ , foils can be exchanged (high Q-values:  $^{150}\text{Nd}$ ,  $^{48}\text{Ca}$ )
- Unique feature: tracking allows to detect  $\beta\beta$ -signature (vertex)
- Demonstrator (= 1st module) currently in commissioning, first data end of 2017
- Design sensitivity:  $T_{1/2} > 10^{26}$  a,  $m_{\beta\beta} \sim 50\text{-}100$  meV

$\beta\beta$ - reconstruction

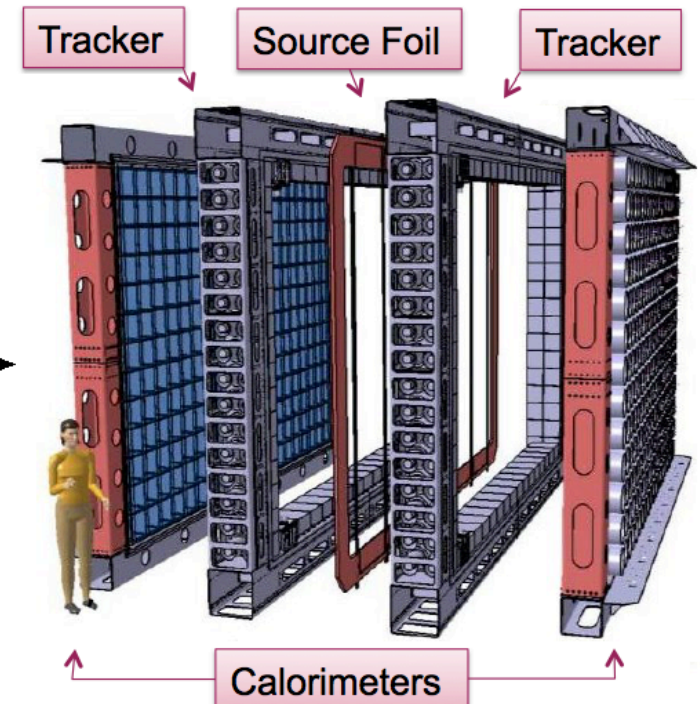


demonstrator module



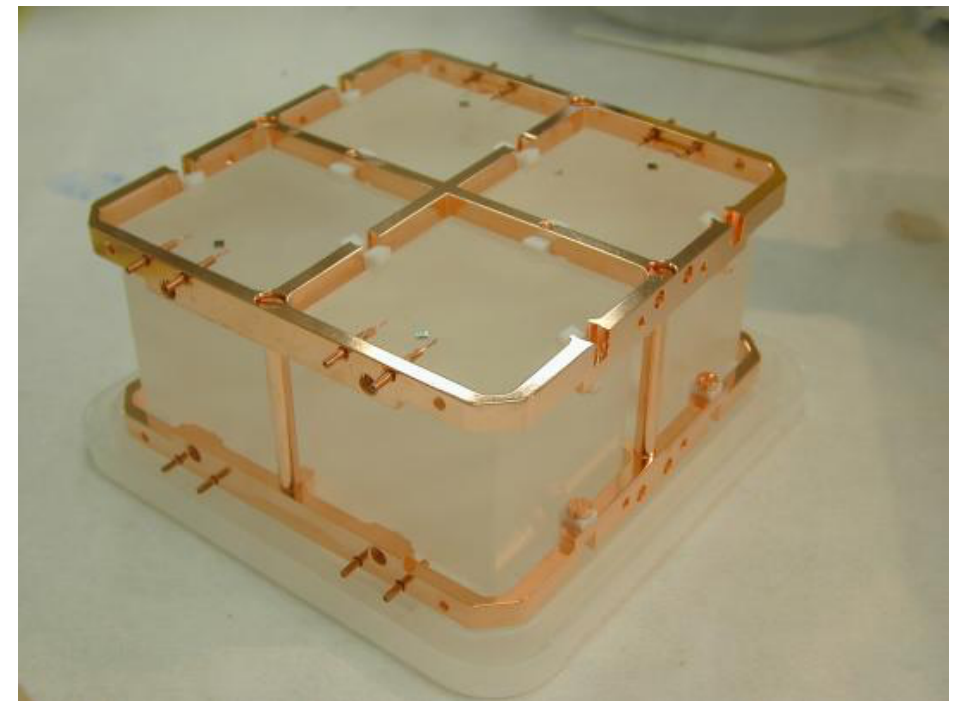
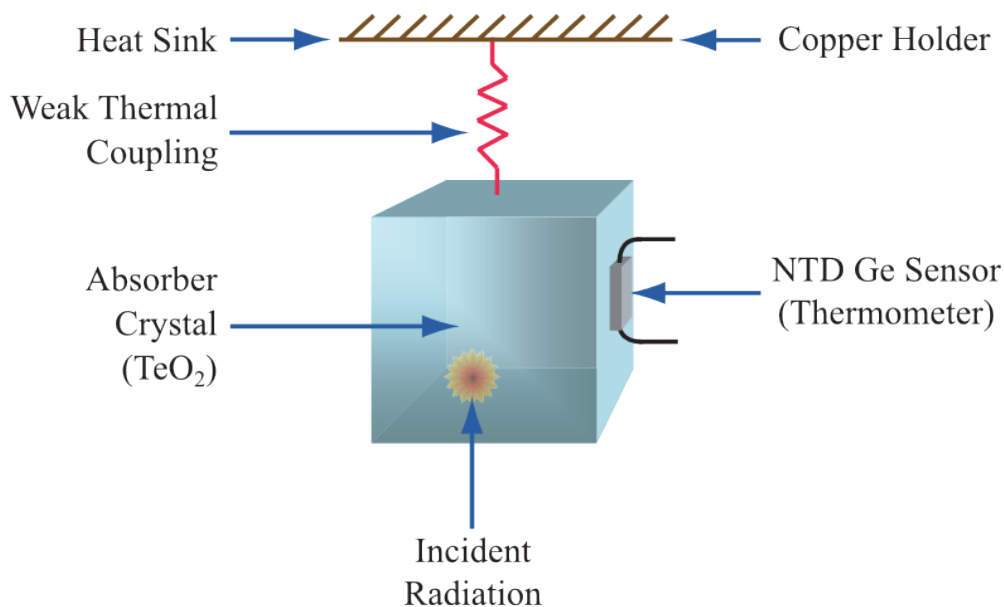
B  $\sim$  2.5 mT  
+ TOF

5-7 kg enriched  
source material



# Cryogenic bolometer technique

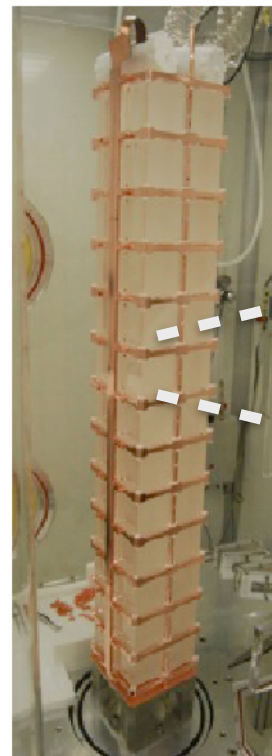
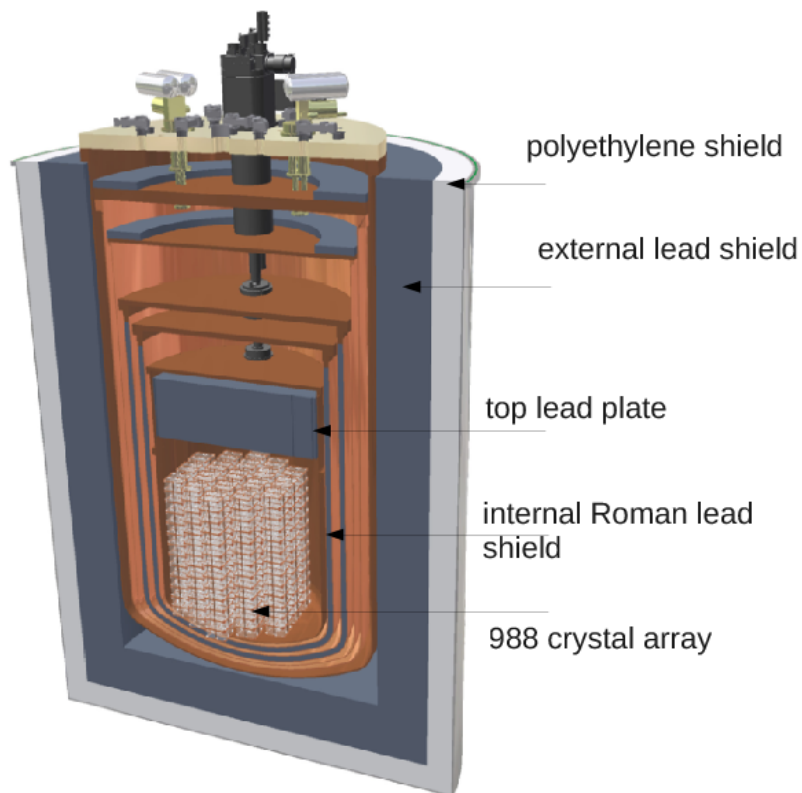
- Electrons create phonons/heat in absorber (e.g.,  $\text{TeO}_2$  crystal)
- Heat capacity:  $\sim (T/T_D)^3$  (Debye Law)
- Example:
  - Operating temperature: 10 mK
  - Temperature change per energy: 10 – 20  $\mu\text{K}/\text{MeV}$
- At  $Q_{\beta\beta} = 2.5 \text{ MeV} \rightarrow \Delta T < 50 \mu\text{K}$



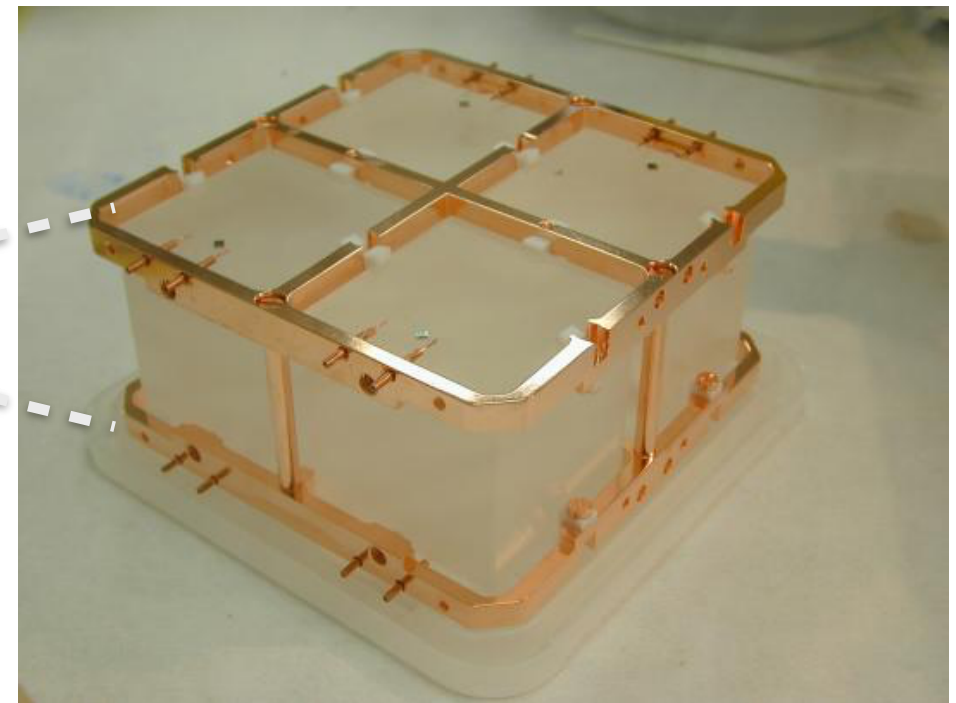


# Cryogenic bolometer: CUORE

- First ton-scale  $0\nu\beta\beta$  exp. with thermal detectors; at Gran Sasso underground laboratory
- $\text{TeO}_2$  detectors & cryo-technology piloted by Cuoricino & CUORE-0 (~40 kg)
- Since Feb. 2017: operation of **988 detectors** at **T ~7 mK**
- Total mass: 742 kg of  $\text{TeO}_2$  → **206 kg of  $^{130}\text{Te}$**



single detector:  $5\times 5\times 5\text{ cm}^3$ , 750 g



19 towers x 13 planes x 4 crystals = 988 crystals

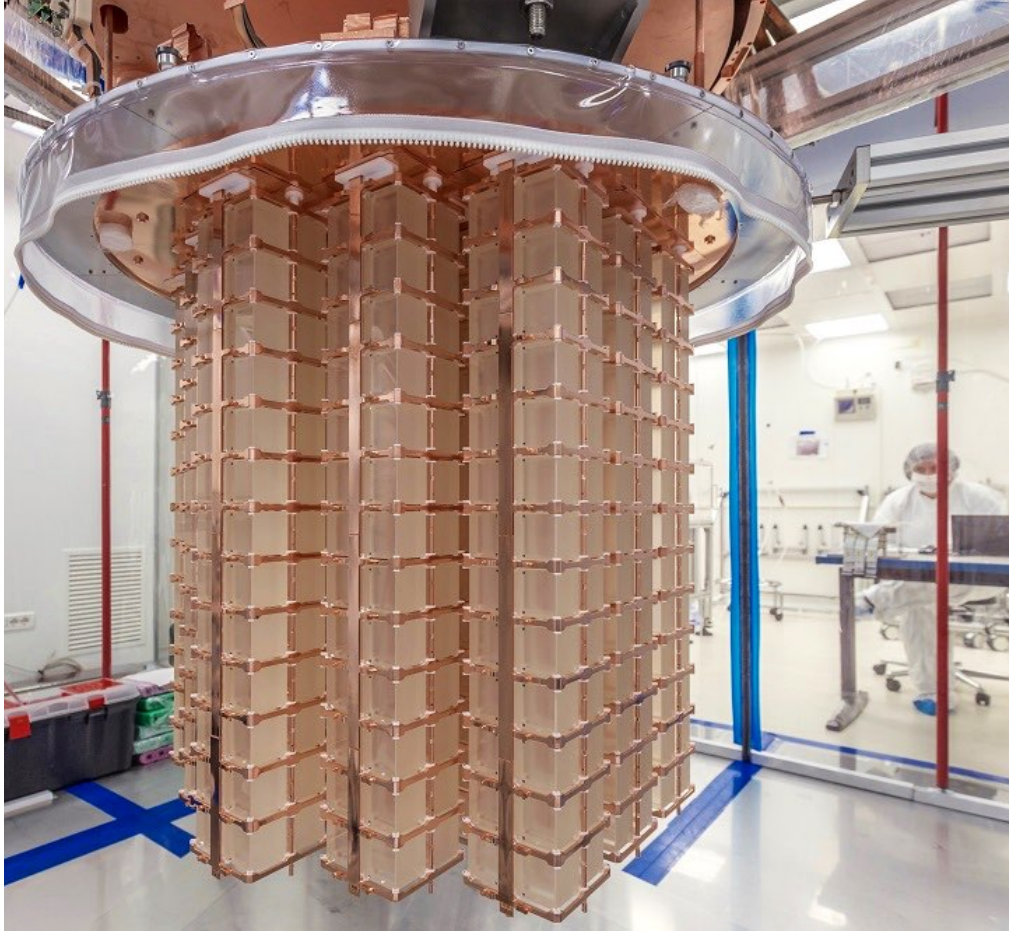


# Cryogenic bolometer: CUORE

The coldest cubic meter in the Universe!

CUORE: at  $\sim 10$  mK

Cosmic microwave background: 2.7 K



19 towers x 13 planes x 4 crystals = 988 crystals



cryogenic platform with helium dilution cryostat

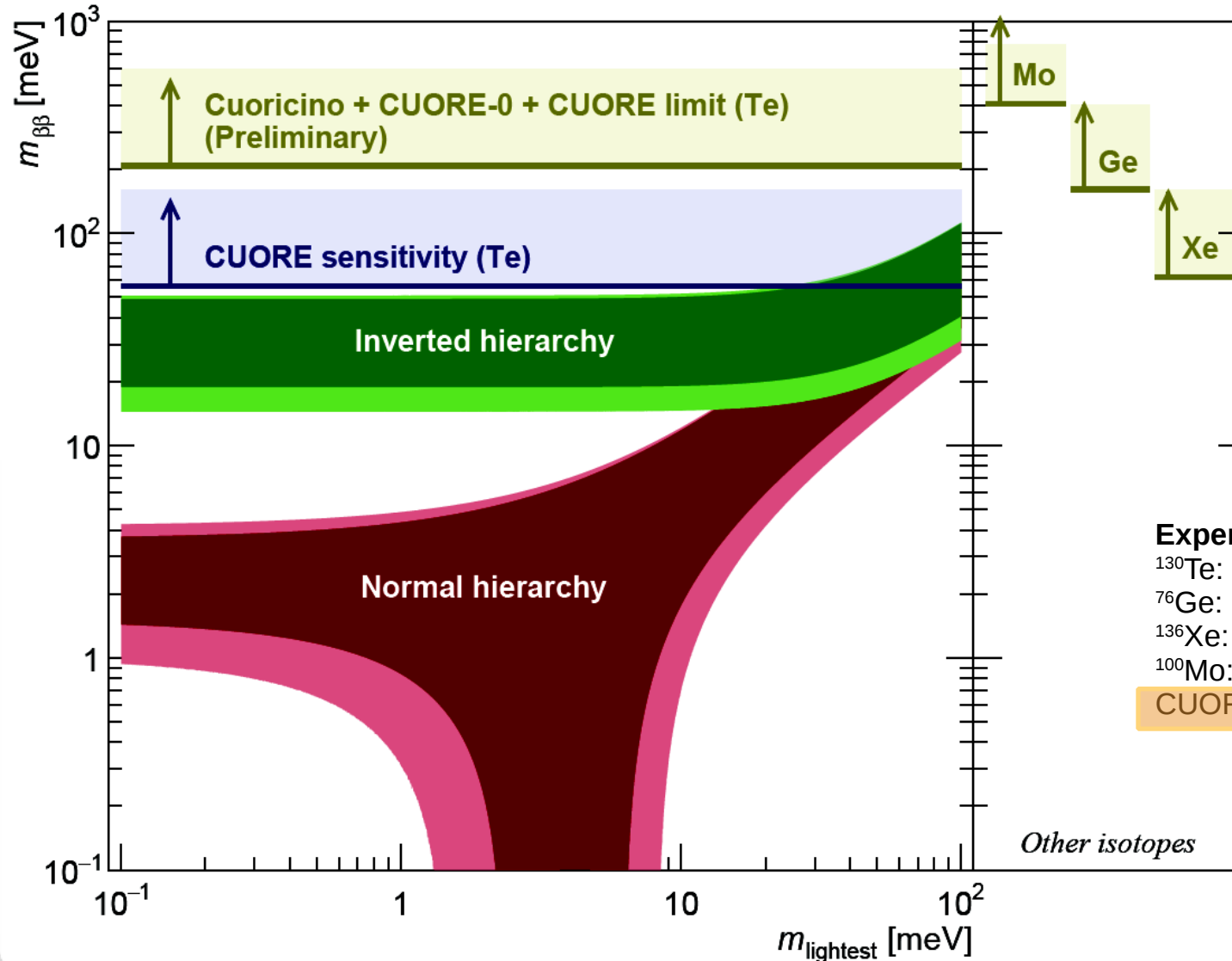
# Cryogenic bolometer: CUORE

First CUORE science run, combined with Cuoricino + CUORE-0:

$$\tau_{1/2}^{0\nu} > 6.6 \times 10^{24} \text{ y (90\% C.L.)}$$



$$m_{\beta\beta} < 210 - 590 \text{ meV}$$



### NME

- Phys. Rev. C 91, 034304 (2015)
- Phys. Rev. C 87, 045501 (2013)
- Phys. Rev. C 91, 024613 (2015)
- Nucl. Phys. A 818, 139 (2009)
- Phys. Rev. Lett. 105, 252503 (2010)

### Experiments

- $^{130}\text{Te}$ :  $6.6 \times 10^{24}$  yr from this analysis
- $^{76}\text{Ge}$ :  $5.3 \times 10^{25}$  yr from Nature 544, 47–52 (2017)
- $^{136}\text{Xe}$ :  $1.1 \times 10^{26}$  yr from Phys. Rev. Lett. 117, 082503 (2016)
- $^{100}\text{Mo}$ :  $1.1 \times 10^{24}$  yr from Phys. Rev. D 89, 111101 (2014)

CUORE sensitivity:  $9.0 \times 10^{25}$  yr

for 5 years of data

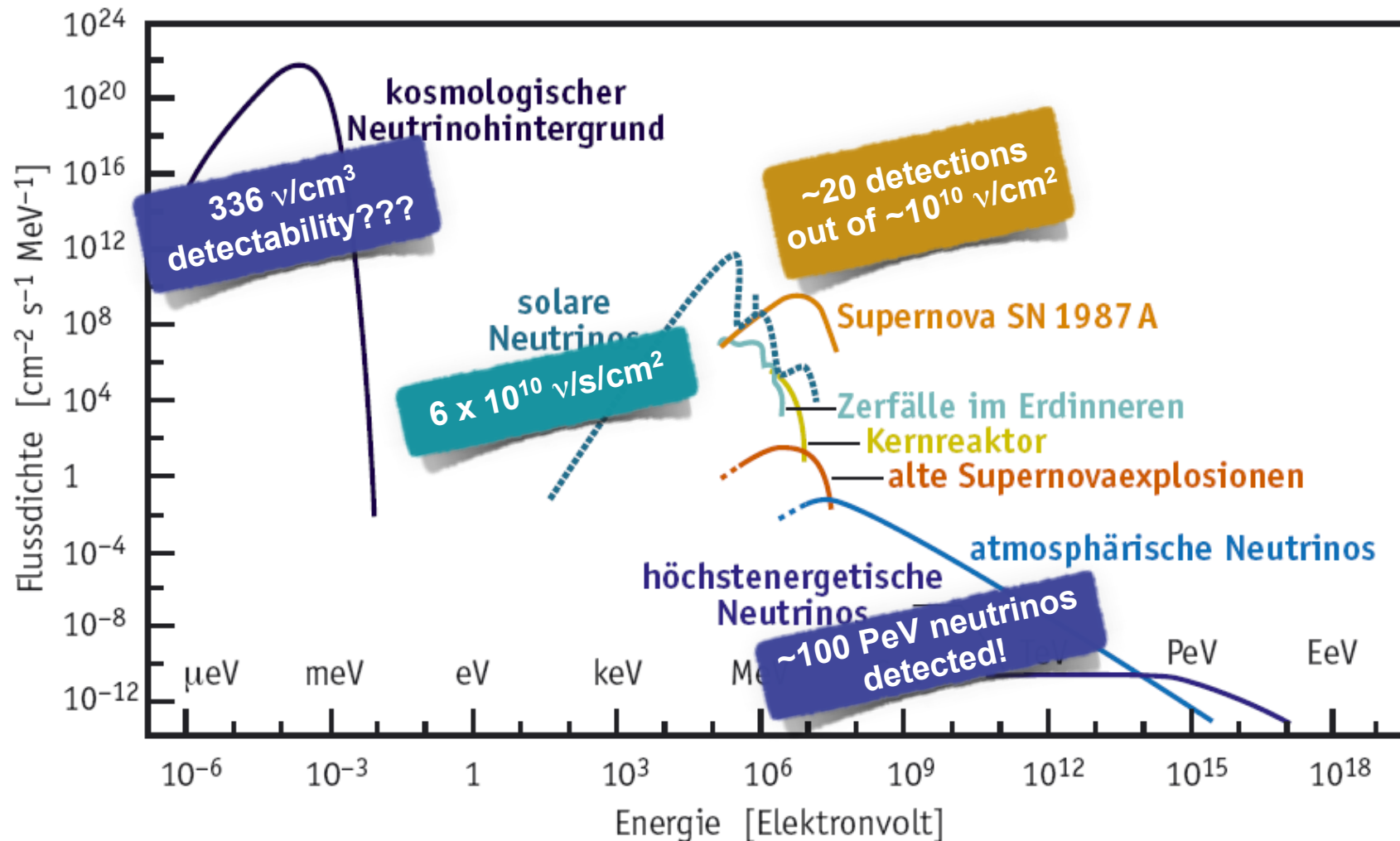
Monica Sisti – Erice, September 18, 2017



# Summary / Take-away (part I)



- We learned a lot about neutrinos since their “invention” in 1930
- We exploit a large variety of neutrino sources in our experiments!



# Summary / Take-away (part I)



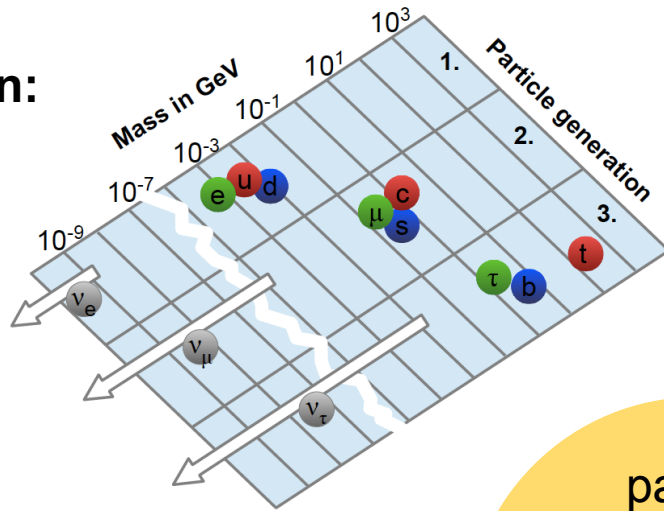
- Massive neutrinos ...
  - are evidence for physics beyond the Standard Model
  - are the only currently known form of Dark Matter  
(their contribution is small, their role not quite fixed yet - what about sterile neutrinos?)
- Neutrinos can point us towards ...
  - novel mass-generating concepts in particle physics  
(open question regarding Dirac or Majorana nature of neutrinos)
  - lepton flavour violation (oscillations) and lepton number violation ( $0\nu\beta\beta$ )
  - leptonic CP violation
- We need to understand the mass pattern of neutrinos  
... and be open for (more) surprises. :)

dedicated experiments  
→ next lecture



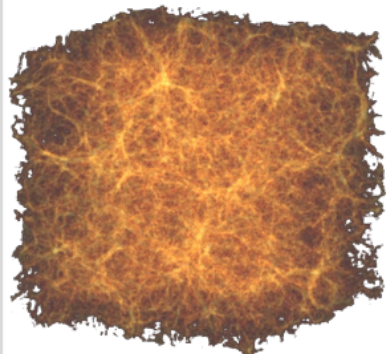
# Massive neutrinos: connecting the micro- and the macro-cosmos

mass generation:  
new concepts

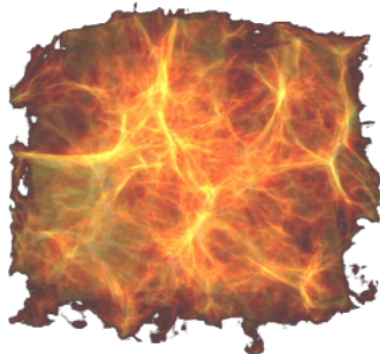


understanding  
astrophysical processes

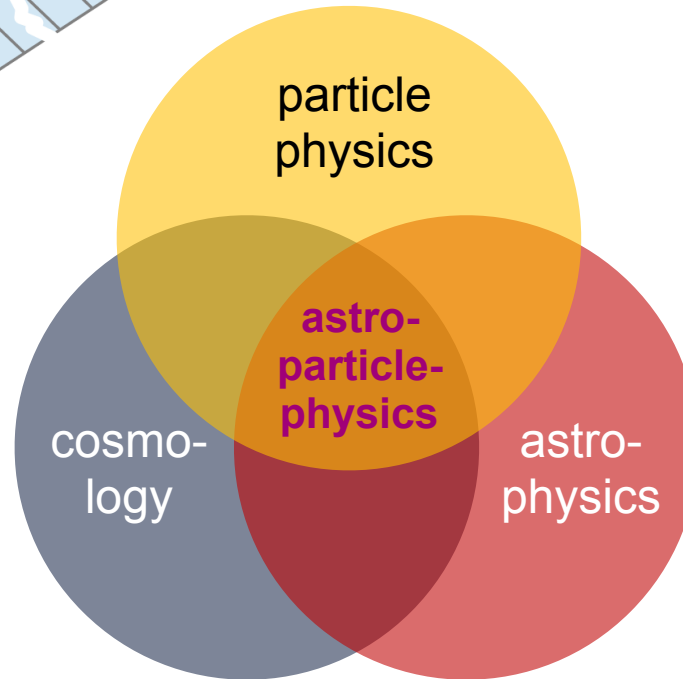
massive neutrinos as  
“cosmic architects”



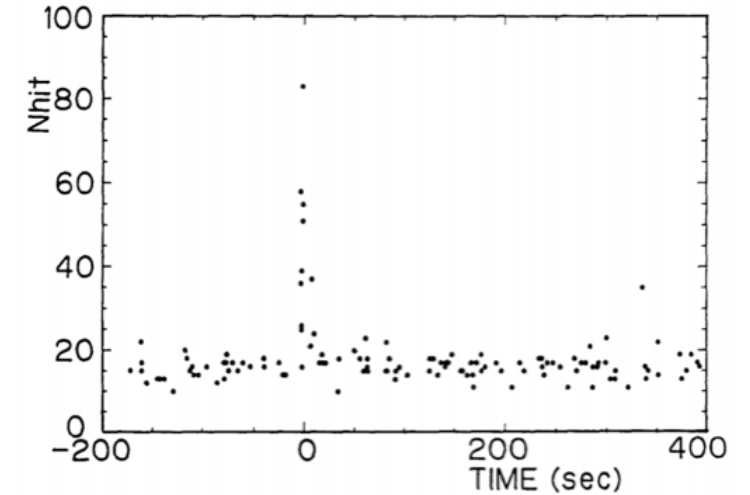
$\Sigma m_\nu = 0 \text{ eV}$



6.9 eV



Neutrino burst from SN 1987a



Matter effects in the sun

